

ON DECAY PROPERTIES OF SOLUTIONS OF THE k -GENERALIZED KDV EQUATION

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ABSTRACT. We prove special decay properties of solutions to the initial value problem associated to the k -generalized Korteweg-de Vries equation. These are related with persistence properties of the solution flow in weighted Sobolev spaces and with sharp unique continuation properties of solutions to this equation. As application of our method we also obtain results concerning the decay behavior of perturbations of the traveling wave solutions as well as results for solutions corresponding to special data.

1. INTRODUCTION

In this work we shall study special decay properties of real solutions to the initial value problem (IVP) associated to the k -generalized Korteweg-de Vries (k -gKdV) equation

$$(1.1) \quad \begin{cases} \partial_t u + \partial_x^3 u + u^k \partial_x u = 0, & x \in \mathbb{R}, \quad k \in \mathbb{Z}^+, \\ u(x, 0) = u_0(x). \end{cases}$$

These decay properties of the solution $u(t)$ will be measured in appropriate weighted $L^2(wdx)$ -spaces.

First, we shall be concerned with asymmetric (increasing) weights, for which the result will be restricted to forward times $t > 0$. In this regard we find the result in [8] for the KdV equation, $k = 1$ in (1.1), in the space $L^2(e^{\beta x} dx)$, $\beta > 0$. There it was shown that the persistence property holds for L^2 -solutions in $L^2(e^{\beta x} dx)$, $\beta > 0$, for $t > 0$ (persistence property in the function space X means that the solution $u(\cdot)$ describes a continuous curve on X , $u \in C([0, T] : X)$). Moreover, formally in this space the operator $\partial_t + \partial_x^3$ becomes $\partial_t + (\partial_x - \beta)^3$ so the solutions of the equation exhibit a parabolic behavior. More precisely, the following result for the KdV equation was proven in [8] (Theorem 11.1 and Theorem 12.1)

Theorem A. *Let $u \in C([0, \infty) : H^2(\mathbb{R}))$ be a solution of the IVP (1.1) with $k = 1$ and*

$$(1.2) \quad u_0 \in H^2(\mathbb{R}) \cap L^2(e^{\beta x} dx), \quad \text{for some } \beta > 0,$$

then

$$(1.3) \quad e^{2\beta x} u \in C([0, \infty) : L^2(\mathbb{R})) \cap C((0, \infty) : H^\infty(\mathbb{R})),$$

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with

$$(1.4) \quad \|u(t)\|_2 = \|u_0\|_2, \quad \|u(t) - u_0\|_{-3,2} \leq Kt, \quad t > 0,$$

$$(1.5) \quad \|e^{\beta x} u(t)\|_2 \leq e^{Kt} \|e^{\beta x} u_0\|_2, \quad t > 0,$$

and

$$(1.6) \quad \int_0^\infty e^{-Kt} \|e^{\beta x} \partial_x u(t)\|_2^2 dt \leq \frac{1}{4\beta} \|e^{\beta x} u_0\|_2^2,$$

where $K = K(\beta, \|u_0\|_2)$.

Moreover, the map data-solution $u_0 \rightarrow u(t)$ is continuous from $L^2(\mathbb{R}) \cap L^2(e^{\beta x} dx)$ to $C([0, T] : L^2(e^{\beta x} dx))$, for any $T > 0$.

On a similar regard, in [4] a unique continuation result was established. This gives an upper bound for the possible space decay of solutions to the IVP (1.1):

Theorem B. *There exists $c_0 > 0$ such that for any pair*

$$u_1, u_2 \in C([0, 1] : H^4(\mathbb{R}) \cap L^2(|x|^2 dx))$$

of solutions of (1.1), if

$$(1.7) \quad u_1(\cdot, 0) - u_2(\cdot, 0), \quad u_1(\cdot, 1) - u_2(\cdot, 1) \in L^2(e^{c_0 x_+^{3/2}} dx),$$

then $u_1 \equiv u_2$.

Above we used the notation: $x_+ = \max\{x; 0\}$. Similarly, we will use later on $x_- = \max\{-x; 0\}$.

In a previous work [14] a similar result for the power 9/4 was established.

The power 3/2 in the exponent in (1.7) reflects the asymptotic behavior of the Airy function. The solution of the initial value problem (IVP)

$$(1.8) \quad \begin{cases} \partial_t v + \partial_x^3 v = 0, \\ v(x, 0) = v_0(x), \end{cases}$$

is given by the group $\{U(t) : t \in \mathbb{R}\}$

$$U(t)v_0(x) = \frac{1}{\sqrt[3]{3t}} Ai\left(\frac{\cdot}{\sqrt[3]{3t}}\right) * v_0(x),$$

where

$$Ai(x) = c \int_{-\infty}^{\infty} e^{ix\xi + i\xi^3/3} d\xi,$$

is the Airy function which satisfies the estimate

$$|Ai(x)| \leq c \frac{e^{-cx_+^{3/2}}}{(1+x_-)^{1/4}}.$$

Observe that Theorem B gives a restriction on the possible decay of a non-trivial solution of (1.1) at two different times. More precisely, taking $u_2 \equiv 0$ one has that if $u_1(t)$ is a solution of the IVP (1.1) such that

$$(1.9) \quad |u_1(x, t)| \leq e^{-a_0 x_+^{3/2}} \quad \text{at } t = 0, 1, \quad \text{for } a_0 \gg 1, \quad \text{then } u_1 \equiv 0.$$

Our first theorem shows that the above result is close to be optimal. By rescaling, the persistence property can not hold in the space $L^2(e^{a_0 x_+^{3/2}} dx)$ in an arbitrary large time interval. However, it does it with a factor $a(t)$ in front of the exponential term $x_+^{3/2}$ which measures how this exponential property decreases with time.

To simplify the exposition, we shall first state our result in the case of the KdV, i.e. $k = 1$ in (1.1) :

Theorem 1.1. *Let a_0 be a positive constant. For any given data*

$$(1.10) \quad u_0 \in L^2(\mathbb{R}) \cap L^2(e^{a_0 x_+^{3/2}} dx),$$

the unique solution of the IVP (1.1) provided by Theorem C below satisfies that for any $T > 0$

$$(1.11) \quad \sup_{t \in [0, T]} \int_{-\infty}^{\infty} e^{a(t)x_+^{3/2}} |u(x, t)|^2 dx \leq C^* = C^*(a_0, \|u_0\|_2, \|e^{a_0 x_+^{3/2}/2} u_0\|_2, T),$$

with

$$(1.12) \quad a(t) = \frac{a_0}{(1 + 27a_0^2 t/4)^{1/2}}.$$

This result extends to the difference of two appropriate solutions of the IVP (1.1) with $k = 1$.

Theorem 1.2. *Let a_0 be a positive constant. Let $u(t), v(t)$ be solutions of the IVP (1.1) with $k = 1$ such that*

$$(1.13) \quad \begin{aligned} u &\in C([0, \infty) : H^1(\mathbb{R}) \cap L^2(|x|dx)) \cap \dots, \\ v &\in C([0, \infty) : H^1(\mathbb{R})) \cap \dots \end{aligned}$$

If

$$(1.14) \quad \int_{-\infty}^{\infty} e^{a_0 x_+^{3/2}} |u(x, 0) - v(x, 0)|^2 dx < \infty,$$

then for any $T > 0$

$$(1.15) \quad \sup_{t \in [0, T]} \int_{-\infty}^{\infty} e^{a(t)x_+^{3/2}} |u(x, t) - v(x, t)|^2 dx \leq C^*,$$

with

$$(1.16) \quad C^* = C^*(a_0, \|u_0\|_{1,2}, \|v_0\|_{1,2}, \| |x|^{1/2} u_0 \|_2, \|u_0 - v_0\|_{1,2}, \|e^{a_0 x_+^{3/2}/2} (u_0 - v_0)\|_2, T),$$

and

$$(1.17) \quad a(t) = \frac{a_0}{(1 + 27a_0^2 t/4)^{1/2}}.$$

The results in Theorem 1.1 and Theorem 1.2 apply to other powers k in the IVP (1.1):

Theorem 1.3. *Let a_0 be a positive constant. For any given data*

$$u_0 \in H^{s_k}(\mathbb{R}) \cap L^2(e^{a_0 x_+^{3/2}} dx)$$

with $s_2 = 1/4$ if $k = 2$, $s_3 = 0$ if $k = 3$, and $s_k > (k - 4)/2k$, $k \geq 4$, then the unique solution of the IVP (1.1) provided by Theorem D below satisfies (1.11) with $a(t)$ as in (1.12) for any $T > 0$ if $k = 2, 3$ and for $T = T(\|u_0\|_{s_k, 2}) > 0$ for $k \geq 4$.

Similarly, for Theorem 1.2.

Remarks:

(a) Following the argument in the proof of Theorem 1.4 in [4] one has that given $a_0 > 0$, $\epsilon > 0$, there exist $u_0 \in \mathcal{S}(\mathbb{R})$, $c_1, c_2 > 0$ and $\Delta T > 0$ such that the corresponding solution $u(x, t)$ of the IVP (1.1) with $k = 1$ satisfies

$$c_1 e^{-(a_0 + \epsilon)x^{3/2}} \leq u(x, t) \leq c_2 e^{-(a_0 - \epsilon)x^{3/2}}, \quad x \gg 1, \quad t \in [0, \Delta T].$$

(b) For different powers k in (1.1) one has that the same statement of Theorem 1.1 is valid except for the last part concerning the continuity of the map data-solution $u_0 \rightarrow u(t)$ which will be continuous from $H^{s_k}(\mathbb{R}) \cap L^2(e^{\beta x} dx)$ to $C([0, T] : L^2(e^{\beta x} dx))$ with s_k and T as in the statement of Theorem 1.3.

(c) For other results involving asymmetric weights of a polynomial type we refer to [11] and [5].

(d) Consider the 1-D semi-linear Schrödinger equation

$$(1.18) \quad \partial_t v = i(\partial_x^2 v + F(v, \bar{v})),$$

with $F : \mathbb{C}^2 \rightarrow \mathbb{C}$, $F \in C^2$ and $F(0) = \partial_u F(0) = \partial_{\bar{u}} F(0) = 0$. As far as we are aware it is unknown whether or not there exist non-trivial solutions $v(t)$ of (1.18) satisfying

$$v(t) \in L^2(e^{a_0 x_+^{1+\epsilon}} dx) \quad t \in [0, T], \quad \text{for some } T > 0, \quad \epsilon > 0.$$

Next, we consider weighted spaces with symmetric weight of the form

$$L^2(\langle x \rangle^b dx) = L^2((1 + x^2)^{b/2} dx)$$

for which persistent properties should hold regardless of the time direction considered, i.e. forward $t > 0$ or backward $t < 0$.

In this setting our second result establishes that for a solution of the IVP (1.1) to satisfy the persistent property in $L^2(\langle x \rangle^b dx)$ it needs to have a similar property in an appropriate Sobolev space $H^s(\mathbb{R})$, i.e. decay in L^2 can only hold if $u(t)$ is regular enough in L^2 :

Theorem 1.4. *Let $u \in C(\mathbb{R} : L^2(\mathbb{R}))$ be the global solution of the IVP (1.1) with $k = 1$ provided by Theorem C below. If there exist $\alpha > 0$ and two different times $t_0, t_1 \in \mathbb{R}$ such that*

$$(1.19) \quad |x|^\alpha u(x, t_0), |x|^\alpha u(x, t_1) \in L^2(\mathbb{R}),$$

then

$$(1.20) \quad u \in C(\mathbb{R} : H^{2\alpha}(\mathbb{R})).$$

Theorem 1.5. *Let $u, v \in C(\mathbb{R} : H^1(\mathbb{R}))$ be global solutions of the IVP (1.1) with $k = 1$ provided by Theorem C. If there exist $\alpha > 1/2$ and two different times $t_0, t_1 \in \mathbb{R}$ such that*

$$(1.21) \quad |x|^\alpha (u(x, t_0) - v(x, t_0)), |x|^\alpha (u(x, t_1) - v(x, t_1)) \in L^2(\mathbb{R}),$$

then

$$(1.22) \quad u - v \in C(\mathbb{R} : H^{2\alpha}(\mathbb{R})).$$

Combining the results in Theorem 1.4 and the results in [13] one obtains:

Corollary 1.1. *Let $s' > s > 0$ and $\beta > 0$. If $u_0 \in H^s(\mathbb{R}) \cap L^2(|x|^\beta dx) - H^{s'}(\mathbb{R})$, then the solution $u = u(x, t)$ of the IVP (1.1) with $k = 1$ provided by Theorem C satisfies:*

(i) $u \in C(\mathbb{R} : H^s(\mathbb{R}) \cap L^2(|x|^\beta dx)$ if $2\beta < s$,

and

(ii) $u(\cdot, t) \notin L^2(|x|^\beta)$ if $2\beta \geq s'$ and $t \neq 0$.

Combining Theorem 1.1 and Theorem 1.4 and taking an initial data $u_0 \in L^2(\mathbb{R})$ with compact support such that

$$u_0 \notin H^s(\mathbb{R}), \quad \forall s > 0,$$

one gets:

Corollary 1.2. *There exists a solution*

$$(1.23) \quad u \in C(\mathbb{R} : L^2(\mathbb{R})) \cap \dots$$

of the IVP (1.1) with $k = 1$ provided by Theorem C such that

$$(1.24) \quad u(\cdot, 0) = u_0(\cdot)$$

has compact support and

$$(1.25) \quad \begin{aligned} u(\cdot, t) &\in C^\infty(\mathbb{R}), & \forall t \neq 0, \\ u(\cdot, t) &\notin L^2(|x|^\epsilon dx), & \forall \epsilon > 0, \forall t \neq 0, \\ u(t) &\in L^2(e^{a(t)x_+^{3/2}/2} dx), & \forall t > 0, \\ u(t) &\in L^2(e^{a(t)x_-^{3/2}/2} dx), & \forall t < 0, \end{aligned}$$

with

$$(1.26) \quad a(t) = \frac{a_0}{(1 + 27a_0^2|t|/4)^{1/2}}.$$

The results in Theorem 1.4, Corollary 1.1 and Corollary 1.2 extend to the other powers $k = 2, 3, 4, \dots$ in (1.1) with the appropriate modifications, accordingly to the value of k , on the regularity and on the length of the time interval $[0, T]$ as it was described in the statement of Theorem 1.3.

The result in Theorem 1.5 holds for any power k in (1.1) and any pair of solutions

$$u, v \in C([-T, T] : H^1(\mathbb{R}))$$

for any $T > 0$ if $k = 2, 3$ and for $T = T(\|u_0\|_{1,2}, \|v_0\|_{1,2})$ for $k \geq 4$.

As a second consequence of our above results one has that there exist compact perturbations of the travelling wave solution with speed c

$$(1.27) \quad u_{k,c}(x, t) = \phi_{k,c}(x - ct),$$

with

$$(1.28) \quad \phi_{k,c}(x) = (c_k c \operatorname{sech}^2(k\sqrt{c}x/2))^{1/k}$$

for the equation in (1.1) which destroy its exponential decay character:

Corollary 1.3. *For a given data of the form*

$$(1.29) \quad u_0(x) = \phi_{k,c}(x) + v_0(x),$$

with $v_0 \in H^1(\mathbb{R})$ compactly supported such that $v_0 \notin H^{1+\epsilon}(\mathbb{R})$ for any $\epsilon > 0$, then the corresponding solution of the IVP provided by Theorem D

$$(1.30) \quad u \in C([-T, T] : H^1(\mathbb{R})) \cap \dots$$

of the IVP (1.1) satisfies that

$$(1.31) \quad u(\cdot, t) \notin L^2(|x|^{1+\epsilon} dx), \quad \forall \epsilon > 0, \quad \forall t \in [-T, T] - \{0\}.$$

Remarks:

(a) As in Corollary 1.2, for $t > 0$ the loss of decay is in the left hand side of \mathbb{R} , and for $t < 0$ in the right hand side of \mathbb{R} .

(b) Combining the results in Theorem 1.4, and its extension for all the equations in (1.1) commented above, with those found in [13] one can also conclude that for $k = 2, 4, 5, \dots$

$$|x|^\alpha u(\cdot, t) \in L^2(\mathbb{R}), \quad \forall t \in [-T, T],$$

and that for $k = 1, 3$ for any $\epsilon > 0$

$$|x|^{\alpha-\epsilon} u(\cdot, t) \in L^2(\mathbb{R}), \quad t \in [-T, T] - [t_0, t_1], \quad |x|^\alpha u(\cdot, t) \in L^2(\mathbb{R}), \quad t \in [t_0, t_1].$$

The main difference between the cases $k = 2, 4, 5, \dots$ and $k = 1, 3$ is that for the later the best available well-posedness results require the use of the spaces $X_{s,b}$ defined in (1.32), which makes fractional weights difficult to handle. In particular, the result in Corollary 1.1 part (i) holds for $2\beta \leq s$ if $k = 2, 4, 5, \dots$

(c) The equivalent of Theorem 1.4 for the semi-linear Schrödinger equation (1.18) in all dimension n was obtained in [12].

We need to recall some results concerning the well-posedness (local and global) of the IVP (1.1). First, we remember the definition of the space $X_{s,b}$ introduced in the context of dispersive equations in [2].

For $s, b \in \mathbb{R}$, $X_{s,b}$ denotes the completion of the Schwartz space $\mathcal{S}(\mathbb{R}^2)$ with respect to the norm

$$(1.32) \quad \|F\|_{X_{s,b}} = \left(\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} (1 + |\tau - \xi^3|)^{2b} (1 + |\xi|)^{2s} |\widehat{F}(\xi, \tau)|^2 d\xi d\tau \right)^{1/2}.$$

The following result was established in [2], see also [10]:

Theorem C. *There exists $b \in (1/2, 1)$ such that for any $u_0 \in L^2(\mathbb{R})$ there exists a unique solution $u(t)$ of the IVP (1.1) with $k = 1$ satisfying*

$$(1.33) \quad u \in C([0, T] : L^2(\mathbb{R})), \quad \text{for any } T > 0,$$

$$(1.34) \quad u \in X_{0,b}, \quad \partial_x(u^2) \in X_{0,b-1}, \quad \partial_t u \in X_{-3,b-1}.$$

Moreover, the map data-solution from $L^2(\mathbb{R})$ into the class defined in (1.33)-(1.34) is Lipschitz for any $T > 0$.

We recall that if $b > 1/2$ one has, using Strichartz estimates and Kato local smoothing effects, that

$$(1.35) \quad \|f\|_{L^4([0,T]:L^\infty(\mathbb{R}))} = \left(\int_0^T \|f(\cdot, t)\|_\infty^4 dt \right)^{1/4} \leq c \|f\|_{X_{0,b}},$$

and

$$(1.36) \quad \|\partial_x f\|_{L^\infty(\mathbb{R}; L^2[0,T])} = \sup_{x \in \mathbb{R}} \left(\int_0^T |\partial_x f(x, t)|^2 dt \right)^{1/2} \leq c \|f\|_{X_{0,b}}.$$

Therefore, combining these estimates and Theorem C one has that the map data-solution is continuous from $L^2(\mathbb{R})$ to $L^4([0, T] : L^\infty(\mathbb{R}))$, for any $T > 0$.

Gathering the local and the global well-posedness results in [9], [3], [15], [7] and [6] one has:

Theorem D. (a) *The IVP (1.1) with $k = 2$ is globally well-posed in $H^s(\mathbb{R})$ for $s \geq 1/4$ (see [3]).*

(b) *The IVP (1.1) with $k = 3$ is locally well-posed in $H^s(\mathbb{R})$ for $s \geq -1/6$, and globally well-posed in $H^s(\mathbb{R})$ for $s > -1/42$ (see [6], [15], [7]).*

(c) *The IVP (1.1) with $k \geq 4$ is locally well-posed in $H^s(\mathbb{R})$ for $s \geq (k - 4)/2k$ (see [9]).*

The rest of this paper is organized as follows: In Section 2 Theorem 1.1 will be proved. Since the proofs of Theorems 1.2 and 1.3 are similar to the proof of Theorem 1.1 we will omit them. The proof of Theorem 1.4 will be given in Section 3. The proof of Theorem 1.5 will be omitted since its proof follows analogous arguments as the ones of Theorem 1.4. Corollaries 1.2 and 1.3 are direct consequence of the previous results.

2. PROOF OF THEOREM 1.1

Consider the IVP

$$(2.1) \quad \begin{cases} \partial_t u + \partial_x^3 u + u \partial_x u = 0, & t \geq 0, \quad x \in \mathbb{R}, \\ u(x, 0) = u_0^\epsilon(x). \end{cases}$$

where

$$(2.2) \quad u_0^\epsilon(x) = \rho_\epsilon * u_0(\cdot + \epsilon)(x) = \int_{-\infty}^{\infty} \frac{1}{\epsilon} \rho\left(\frac{y}{\epsilon}\right) u_0(x + \epsilon - y) dy,$$

where $\rho \in C^\infty(\mathbb{R})$ with $\text{supp } \rho \subset \{x \in \mathbb{R} : |x| \leq 1\}$, $\rho \geq 0$ and $\int \rho(x) dx = 1$. We claim that for any $\epsilon \in (0, 1)$

$$(2.3) \quad u_0^\epsilon \in H^\infty(\mathbb{R}) \cap L^2(e^{a_0 x_+^{3/2}} dx).$$

To see (2.3) we use Minkowski's integral inequality to write:

$$(2.4) \quad \begin{aligned} \|e^{a_0 x_+^{3/2}/2} u_0^\epsilon\|_2 &= \left(\int e^{a_0 x_+^{3/2}} \left| \int \frac{1}{\epsilon} \rho\left(\frac{y}{\epsilon}\right) u_0(x + \epsilon - y) dy \right|^2 dx \right)^{1/2} \\ &= \left(\int \left| \int e^{a_0 x_+^{3/2}/2} \rho_\epsilon(y) u_0(x + \epsilon - y) dy \right|^2 dx \right)^{1/2} \\ &\leq \int \rho_\epsilon(y) \left(\int e^{a_0 x_+^{3/2}} |u_0(x + \epsilon - y)|^2 dx \right)^{1/2} dy \\ &\leq \int \rho_\epsilon(y) \left(\int e^{a_0(x-\epsilon+y)_+^{3/2}} |u_0(x)|^2 dx \right)^{1/2} dy \\ &\leq \int \rho_\epsilon(y) \left(\int e^{a_0 x_+^{3/2}} |u_0(x)|^2 dx \right)^{1/2} dy = \left(\int e^{a_0 x_+^{3/2}} |u_0(x)|^2 dx \right)^{1/2}, \end{aligned}$$

since for $y \in \text{supp}(\rho_\epsilon)$ one has that $-\epsilon + y \leq 0$.

Also we recall that

$$(2.5) \quad \lim_{\epsilon \downarrow 0} \|u_0^\epsilon - u_0\|_2 = 0.$$

Therefore, for any $\epsilon \in (0, 1)$ the corresponding solution $u^\epsilon(\cdot)$ of the IVP (2.1) has the properties stated in Theorems A and C.

Next, for any $N \in \mathbb{Z}^+$, $N \gg 1$ we define the weight

$$(2.6) \quad \varphi_N(x, t) = \begin{cases} e^{a(t)/4}, & x \leq 0, \\ e^{a(t)\theta(x)}, & 0 \leq x \leq 1, \\ e^{a(t)x^{3/2}}, & 1 \leq x \leq N, \\ P_2(x, t), & x \geq N, \end{cases}$$

where

$$(2.7) \quad a'(t) + \frac{27}{8}a^3(t) = 0, \quad a(0) = a_0,$$

i.e.

$$(2.8) \quad a(t) = \frac{a_0}{(1 + 27a_0^2 t/4)^{1/2}} \in (0, a_0], \quad \forall t \geq 0,$$

$$(2.9) \quad \theta(x) = \frac{1}{4} + \frac{15}{8}x^3 - \frac{12}{8}x^4 + \frac{3}{8}x^5,$$

and

$$(2.10) \quad \begin{aligned} P_2(x, t) &= e^{a(t)N^{3/2}} + \frac{3}{2}a(t)N^{1/2}e^{a(t)N^{3/2}}(x - N) \\ &+ \left(\left(\frac{3}{2}a(t)N^{1/2}\right)^2 + \frac{3}{4}a(t)N^{-1/2}\right)e^{a(t)N^{3/2}}\frac{(x - N)^2}{2}. \end{aligned}$$

Remarks :

(i) For any $N \in \mathbb{Z}^+$, $x \geq 0$ and $t \geq 0$ one has

$$(2.11) \quad \varphi_N(x, t) \leq C(a_0) e^{a(t)x^{3/2}},$$

(ii) the function $\theta(\cdot)$ matches the values of $x^{3/2}$ at $x = 1$ and those ones of the constant function $f(x) = 1/4$ at $x = 0$ and their derivatives up to order two and $P_2(x, t)$ matches those ones of $e^{a(t)x^{3/2}}$ at $x = N$ up to order two. Hence, $\varphi_N(\cdot, t) \in C^2(\mathbb{R})$ and $\varphi_N(\cdot, t) \in C^3(\mathbb{R} - \{0, 1, N\})$ for all $t \geq 0$,

(iii) one has that

$$\theta''(x) = \frac{3x}{4}(15 - 24x + 10x^2) = \frac{3x}{4}\left(\left(\sqrt{10}x - \frac{12}{\sqrt{10}}\right)^2 + \frac{3}{5}\right) \geq 0,$$

with $\theta'(0) = 0$, hence $\theta'(x) \geq 0$, $x \in (0, 1]$.

(iv) $P_2(x, t) \geq 0$ and

$$\begin{aligned} \partial_x P_2(x, t) &= \frac{3}{2}a(t)N^{1/2}e^{a(t)N^{3/2}} \\ &+ \left(\left(\frac{3}{2}a(t)N^{1/2}\right)^2 + \frac{3}{4}a(t)N^{-1/2}\right)e^{a(t)N^{3/2}}(x - N) \geq 0, \end{aligned}$$

for $x \geq N$ and $t \geq 0$.

Therefore,

$$(2.12) \quad \varphi_N(x, t), \quad \partial_x \varphi_N(x, t) \geq 0, \quad (x, t) \in \mathbb{R} \times [0, \infty).$$

Moreover, for $x \geq N$

$$\begin{aligned}
(2.13) \quad \partial_x P_2(x, t) &\leq \frac{3}{2} a N^{1/2} P_2(x, t) + \frac{3}{2} a N^{1/2} e^{aN^{3/2}} (x - N) \left(\frac{3}{2} a N^{1/2} + \frac{1}{2N} \right) \\
&\leq \frac{3}{2} a N^{1/2} P_2(x, t) + \left(\frac{3}{2} a N^{1/2} + \frac{1}{2N} \right) P_2(x, t) \\
&\leq (1 + 3a(t)N^{1/2}) P_2(x, t) \\
&\leq (1 + 3a_0 x^{1/2}) \varphi_N(x, t).
\end{aligned}$$

Next, we multiply the equation in (2.1) by $u^\epsilon \varphi_N(x, t)$ where u^ϵ denotes the solution of the IVP (2.1), integrating the result and formally using integration by parts one has

$$\begin{aligned}
(2.14) \quad \frac{d}{dt} \int (u^\epsilon)^2 \varphi_N dx + 3 \int (\partial_x u^\epsilon)^2 \partial_x \varphi_N dx \\
= \int (u^\epsilon)^2 (\partial_x^3 \varphi_N + \partial_t \varphi_N) dx + \frac{2}{3} \int (u^\epsilon)^3 \partial_x \varphi_N dx.
\end{aligned}$$

To justify the integration by parts used to obtain (2.14) we observe that by Theorems A and Theorem C

$$(2.15) \quad u^\epsilon \in C([0, \infty) : H^\infty(\mathbb{R})), \quad e^{\beta x} u^\epsilon \in C([0, \infty) : L^2(\mathbb{R})) \quad \forall \beta > 0.$$

In general, if for some $\beta > 0$ $e^{\beta x} f$, $\partial_x^2 f \in L^2(\mathbb{R})$, then $e^{\beta x/2} \partial_x f \in L^2(\mathbb{R})$, and so $e^{\beta x/2} f \in H^1(\mathbb{R})$ (in particular, $e^{\beta x/2} f(x) \rightarrow 0$ as $|x| \rightarrow \infty$) since

$$(2.16) \quad \int e^{\beta x} (\partial_x f)^2 dx \leq \beta^2 \int e^{\beta x} f^2 dx + \left| \int e^{\beta x} f \partial_x^2 f dx \right|.$$

To prove (2.16) one first assumes that $f \in H^2(\mathbb{R})$ with compact support to obtain (2.16) by integration by parts, and then use the density of this class to get the desired result.

Thus, reapplying the last argument and using (2.15) it follows that for

$$e^{\beta x} \partial_x^j u^\epsilon \in L^2(\mathbb{R}), \quad j = 1, 2, 3,$$

and so for any $t \geq 0$

$$(2.17) \quad e^{\beta x} \partial_x^j u^\epsilon(x, t) \rightarrow 0 \quad j = 0, 1, 2, \quad \text{as } |x| \rightarrow \infty,$$

which justify all integration by parts used to get (2.14) since at $+\infty$, $\varphi_N(x, t)$ as a function of x is a polynomial of order two.

Since

$$\partial_x \varphi_N(x, t) \geq 0,$$

one can omit the second term on the left hand side (l.h.s.) of (2.14) to write

$$(2.18) \quad \frac{d}{dt} \int (u^\epsilon)^2 \varphi_N dx \leq \int (u^\epsilon)^2 (\partial_x^3 \varphi_N + \partial_t \varphi_N) dx + \frac{2}{3} \int (u^\epsilon)^3 \partial_x \varphi_N dx.$$

We shall consider the right hand side (r.h.s) of (2.18) in four different domains:

$$(2.19) \quad \text{(a) } x \leq 0, \quad \text{(b) } 0 \leq x \leq 1, \quad \text{(c) } 1 \leq x \leq N, \quad \text{(d) } x \geq N.$$

In the domain (a) ($x \leq 0$) one has

$$\partial_x^j \varphi_N(x, t) = 0 \quad j = 1, 2, 3, \quad \text{and} \quad \partial_t \varphi_N(x, t) \leq 0,$$

therefore the contribution of this domain to the r.h.s. of (2.18) is non-positive so it does not have to be considered.

In the domain (b) ($0 \leq x < 1$) one has

$$\begin{aligned} \varphi_N(x, t) &\leq \varphi_N(1, t) = e^{a(t)} \leq e^{a_0}, \\ \partial_t \varphi_N(x, t) &= a'(t) \varphi_N(x, t) \leq 0, \\ \partial_x \varphi_N(x, t) &= a(t) \theta'(x) \varphi_N(x, t) \leq c a_0 \varphi_N(x, t), \\ \partial_x^2 \varphi_N(x, t) &= (a(t) \theta''(x) + (a(t) \theta'(x))^2) \varphi_N(x, t) \leq c (a_0 + a_0^2) \varphi_N(x, t), \\ \partial_x^3 \varphi_N(x, t) &= (a(t) \theta^{(3)}(x) + 3(a(t))^2 \theta''(x) \theta'(x) + (a(t) \theta'(x))^{(3)}) \varphi_N(x, t) \\ &\leq c (a_0 + a_0^2 + a_0^3) \varphi_N(x, t). \end{aligned}$$

Therefore

$$(2.20) \quad \int_0^1 (u^\epsilon)^2 (\partial_x^3 \varphi_N + \partial_t \varphi_N) dx \leq K(a_0) \int_0^1 (u^\epsilon)^2 \varphi_N dx,$$

and

$$(2.21) \quad \int_0^1 (u^\epsilon)^3 \partial_x \varphi_N dx \leq K(a_0) \|u^\epsilon(t)\|_{L^\infty(0 \leq x \leq 1)} \int_0^1 (u^\epsilon)^2 \varphi_N dx.$$

In the domain (c) ($1 \leq x < N$) one has

$$(2.22) \quad \partial_x^3 \varphi_N + \partial_t \varphi_N = \left(\frac{27}{8} a^3 x^{3/2} + \frac{27}{8} a^2 - \frac{3}{8} a x^{-3/2} + a' x^{3/2} \right) \varphi_N(x, t).$$

For the third term on the r.h.s. of (2.22) we observe that

$$(2.23) \quad -\frac{3}{8} a(t) x^{-3/2} e^{a(t) x^{3/2}} \leq 0,$$

so it can be omitted. Next, we use that

$$(2.24) \quad a'(t) + \frac{27}{8} a^3(t) = 0,$$

which combined with (2.23) and (2.24) gives us the estimate

$$(2.25) \quad \int_1^N (u^\epsilon)^2 (\partial_x^3 \varphi_N + \partial_t \varphi_N) dx \leq \frac{27}{8} a^2(t) \int_1^N (u^\epsilon)^2 \varphi_N dx \leq \frac{27}{8} a_0^2 \int_1^N (u^\epsilon)^2 \varphi_N dx.$$

For the other term on the r.h.s. of (2.18) in this domain one sees that

$$(2.26) \quad \begin{aligned} \int_1^N (u^\epsilon)^3 \partial_x \varphi_N dx &= \frac{3}{2} \int_1^N a(t) x^{1/2} (u^\epsilon)^3 \varphi_N dx \\ &\leq \frac{3}{2} a_0 \|x^{1/2} u^\epsilon(t)\|_{L^\infty(1 < x < N)} \int_1^N (u^\epsilon)^2 \varphi_N dx. \end{aligned}$$

Finally, we consider the domain (d) ($x \geq N$). Since in this domain

$$\partial_x^3 P_2(x, t) \equiv 0, \quad \partial_t P_2(x, t) = a'(t) \{\cdot\} \leq 0,$$

one has

$$\int_N^\infty (u^\epsilon)^2 (\partial_x^3 \varphi_N + \partial_t \varphi_N) dx \leq 0,$$

so the contribution of this term on the r.h.s. of (2.18) does not need to be considered. To bound the contribution of the other term on the r.h.s. of (2.18) using (2.13) we write

$$(2.27) \quad \begin{aligned} \int_N^\infty (u^\epsilon)^3 \partial_x \varphi_N dx &= \int_N^\infty (u^\epsilon)^3 (1 + 3a_0 x^{1/2}) P_2(x, t) dx \\ &\leq K(a_0) \|x^{1/2} u^\epsilon(t)\|_{L^\infty(x \geq N)} \int_N^\infty (u^\epsilon)^2 \varphi_N(x, t) dx. \end{aligned}$$

It remains to estimate the terms in the L^∞ -norm in (2.21), (2.26) and (2.27). We estimate the terms (2.26) and (2.27), the estimation in (2.21) being similar. For (2.26) and (2.27), using (1.5) in Theorem A and an estimate similar to that in (2.4), we have

$$(2.28) \quad \begin{aligned} \|x^{1/2} u^\epsilon(t)\|_{L^\infty(x > 1)} &\leq \|e^x u^\epsilon(t)\|_\infty \leq \|e^x u^\epsilon(t)\|_2^{1/2} \|\partial_x(e^x u^\epsilon(t))\|_2^{1/2} \\ &\leq \|e^x u^\epsilon(t)\|_2^{1/2} (\|e^x u^\epsilon(t)\|_2^{1/2} + \|e^x \partial_x u^\epsilon(t)\|_2^{1/2}) \\ &\leq 2 \|e^x u^\epsilon(t)\|_2 + \|e^x \partial_x u^\epsilon(t)\|_2 \\ &\leq 2e^M \|e^x u_0^\epsilon\|_2 + \|e^x \partial_x u^\epsilon(t)\|_2 \\ &\leq 2e^M \|e^x u_0\|_2 + \|e^x \partial_x u^\epsilon(t)\|_2. \end{aligned}$$

From (1.6) in Theorem A and an argument similar to that in (2.4) we obtain the bound for the integral in time interval $[0, T]$ of the last term in (2.28)

$$(2.29) \quad \begin{aligned} \int_0^T \|e^x \partial_x u^\epsilon(t)\|_2 dt &\leq T^{1/2} \left(\int_0^T \|e^x \partial_x u^\epsilon(t)\|_2^2 dt \right)^{1/2} \\ &\leq T^{1/2} \frac{1}{4} e^{MT} \|e^x u_0^\epsilon\|_2^2 \leq T^{1/2} \frac{1}{4} e^{MT} \|e^x u_0\|_2^2, \end{aligned}$$

with $M = M(\|u_0\|_2)$.

Inserting the above estimates in (2.18), using Gronwall's inequality, and applying (2.11) with $t = 0$, one gets that

$$(2.30) \quad \sup_{t \in [0, T]} \int_{-\infty}^{\infty} |u^\epsilon(x, t)|^2 \varphi_N(x, t) dx \leq C^* = C^*(a_0, \|u_0\|_2, \|e^{a_0 x_+^{3/2}/2} u_0\|_2, T).$$

We observe that for each $N \in \mathbb{Z}^+$ fixed there exists c_N such that

$$(2.31) \quad \varphi_N(x, t) \leq c_N e^{a_0 x}, \quad x \geq 0,$$

and by the continuity of the map data-solution, in Theorem A it follows that

$$(2.32) \quad \sup_{t \in [0, T]} \int_{-\infty}^{\infty} e^{a_0 x} |u^\epsilon(x, t) - u(x, t)|^2 dx \rightarrow 0 \quad \text{as } \epsilon \downarrow 0.$$

Combining (2.31) and (2.32) one concludes that

$$(2.33) \quad \sup_{t \in [0, T]} \int_{-\infty}^{\infty} \varphi_N(x, t) |u^\epsilon(x, t) - u(x, t)|^2 dx \rightarrow 0 \quad \text{as } \epsilon \downarrow 0.$$

Hence, for any $t \in [0, T]$

$$\int_{-\infty}^{\infty} \varphi_N(x, t) |u^\epsilon(x, t)|^2 dx \rightarrow \int_{-\infty}^{\infty} \varphi_N(x, t) |u(x, t)|^2 dx \quad \text{as } \epsilon \downarrow 0,$$

and consequently, for any $t \in [0, T]$

$$(2.34) \quad \int_{-\infty}^{\infty} \varphi_N(x, t) |u(x, t)|^2 dx \leq C^* = C^*(a_0, \|u_0\|_2, \|e^{a_0 x_+^{3/2}/2} u_0\|_2, T).$$

Finally, letting $N \uparrow \infty$ in (2.34) Fatou's lemma yields the desired result

$$(2.35) \quad \int_{-\infty}^{\infty} e^{a(t) x_+^{3/2}} |u(x, t)|^2 dx \leq C^* = C^*(a_0, \|u_0\|_2, \|e^{a_0 x_+^{3/2}/2} u_0\|_2, T).$$

3. PROOF OF THEOREM 1.4

Without loss of generality we assume $t_0 = 0$ and $t_1 > 0$.

First we shall consider the case $\alpha \in (0, 1/2]$.

For $x \geq 0$, $N \in \mathbb{Z}^+$ and $\alpha > 0$ we define

$$(3.1) \quad \varphi_{N, \alpha}(x) = \begin{cases} (1 + x^4)^{\alpha/2} - 1, & x \in [0, N], \\ (2N)^{2\alpha}, & x \geq 10N, \end{cases}$$

with $\varphi_{N, \alpha} \in C^3((0, \infty))$, $\varphi_{N, \alpha}(x) \geq 0$, $\varphi'_{N, \alpha} \geq 0$ and for $\alpha \in (0, 1/2]$

$$(3.2) \quad |\varphi_{N, \alpha}^{(j)}(x)| \leq c, \quad j = 1, 2, 3, \quad c \text{ independent of } N.$$

Let $\phi_{N,\alpha}$ be the odd extension of $\varphi_{N,\alpha}$, i.e.

$$(3.3) \quad \phi_{N,\alpha}(x) = \phi_N(x) = \begin{cases} \varphi_{N,\alpha}(x), & x \geq 0, \\ -\varphi_{N,\alpha}(-x), & x \leq 0. \end{cases}$$

Notice that

$$(3.4) \quad \phi'_N(x) \geq 0, \quad \forall x \in \mathbb{R}, \quad \phi_N \in C^3(\mathbb{R}), \quad \|\phi_N\|_\infty = (2N)^{2\alpha}.$$

Next, we consider a sequence of data $(u_{0,m})_{m \in \mathbb{Z}^+} \subset \mathcal{S}(\mathbb{R})$ such that

$$(3.5) \quad \|u_0 - u_{0,m}\|_2 \rightarrow 0 \quad \text{as } m \uparrow \infty,$$

and denote by u_m the solution of the IVP (1.1) with $k = 1$ and data $u_m(x, 0) = u_{0,m}(x)$. From the results in [8] one has that

$$(3.6) \quad u_m \in C(\mathbb{R} : \mathcal{S}(\mathbb{R})).$$

By the continuous dependence of the solution upon the data (see Theorem C and comments after its statement) one has that for each $T > 0$

$$(3.7) \quad \begin{aligned} (a) \quad & \sup_{t \in [0, T]} \|u(t) - u_m(t)\|_2 \rightarrow 0 \quad \text{as } m \uparrow \infty, \\ (b) \quad & \int_{-T}^T \|u(t) - u_m(t)\|_\infty^4 dt \rightarrow 0 \quad \text{as } m \uparrow \infty, \\ (c) \quad & \sup_{x \in \mathbb{R}} \int_{-T}^T |\partial_x(u - u_m)(x, t)|^2 dt \rightarrow 0 \quad \text{as } m \uparrow \infty. \end{aligned}$$

Since $u_m \in C(\mathbb{R} : \mathcal{S}(\mathbb{R}))$ satisfies the equation in (1.1) with $k = 1$, multiplying it by $u_m \phi_N$ after integration by parts (justified since ϕ_N is bounded) one gets

$$(3.8) \quad \frac{d}{dt} \int (u_m)^2 \phi_N dx + 3 \int (\partial_x u_m)^2 \phi'_N dx = \int (u_m)^2 \phi_N^{(3)} dx + \frac{2}{3} \int (u_m)^3 \phi'_N dx.$$

We observe that for m large enough

$$(3.9) \quad \begin{aligned} & \left| \int (u_m)^2 \phi_N^{(3)} dx \right| \leq c \|u_{0,m}\|_2^2 \leq 2c \|u_0\|_2^2, \\ & \left| \int (u_m)^3 \phi'_N dx \right| \leq c \|u_m(t)\|_\infty \|u_{0,m}\|_2^2 \leq 2c \|u_m(t)\|_\infty \|u_0\|_2^2. \end{aligned}$$

Integrating in the time interval $[0, t_1]$ the identity (3.8) and using (3.9) it follows that

$$(3.10) \quad \begin{aligned} \int_0^{t_1} \int (\partial_x u_m)^2(x, t) \phi'_N(x) dx dt & \leq \|(u_m(t_1))^2 \phi_N\|_1 + \|(u_{0,m})^2 \phi_N\|_1 \\ & + c t_1 \|u_0\|_2^2 + c t_1^{3/4} \|u_0\|_2^2 \left(\int_0^{t_1} \|u_m(t)\|_\infty^4 dt \right)^{1/4}, \end{aligned}$$

where c denotes a constant whose value may change from line to line and is independent of the initial parameters of the problem. Letting $m \uparrow \infty$, using (3.7) part (c), Theorem C, and (1.35), one gets that

$$(3.11) \quad \overline{\lim}_{m \uparrow \infty} \int_0^{t_1} \int (\partial_x u_m)^2(x, t) \phi'_N(x) dx dt \leq \|(u(t_1))^2 \phi_N\|_1 + \|(u_0)^2 \phi_N\|_1 \\ + c t_1 \|u_0\|_2^2 + c t_1^{3/4} \|u_0\|_2^2 \left(\int_0^{t_1} \|u(t)\|_\infty^4 dt \right)^{1/4} \leq M,$$

with $M = M(\|\langle x \rangle^\alpha u_0\|_2, \|\langle x \rangle^\alpha u(t_1)\|_2)$. Next, we use (3.7) part (c) to conclude that for any $\bar{N} \in \mathbb{Z}^+$ fixed

$$(3.12) \quad \partial_x u_m \rightarrow \partial_x u \quad \text{in } L^2([-T, T] \times [-\bar{N}, \bar{N}]) \quad \text{as } m \uparrow \infty.$$

Since ϕ'_N has compact support one gets that

$$(3.13) \quad \int_0^{t_1} \int (\partial_x u)^2(x, t) \phi'_N(x) dx dt \leq M.$$

Finally, we recall that $\phi'_N(x)$ is even, $\phi'_N(x) \geq 0$, and for $x > 1$

$$\phi'_N(x) \rightarrow \frac{2\alpha x^3}{(1+x^4)^{1-\alpha/2}} \sim \langle x \rangle^{2\alpha-1},$$

therefore using Fatou's lemma in (3.13) one concludes that

$$(3.14) \quad \int_0^{t_1} \int_{|x| \geq 1} (\partial_x u)^2(x, t) \langle x \rangle^{2\alpha-1} dx dt \leq M.$$

Since (see Theorem C and (1.36))

$$\int_0^{t_1} \int_{|x| \leq 1} (\partial_x u)^2(x, t) dx dt \leq M,$$

one concludes that

$$(3.15) \quad \int_0^{t_1} \int (\partial_x u)^2(x, t) \langle x \rangle^{2\alpha-1} dx dt \leq M.$$

Once we have obtained (3.15) we reapply the above argument with $\psi_{N,\alpha}(x) = \psi_N(x)$ the even extension of $\varphi_{N,\alpha}$ instead of $\phi_N(x)$, i.e.

$$(3.16) \quad \psi_{N,\alpha}(x) = \psi_N(x) = \begin{cases} \varphi_{N,\alpha}(x), & x \geq 0, \\ \varphi_{N,\alpha}(-x), & x \leq 0. \end{cases}$$

We observe that

$$(3.17) \quad |\psi'_N(x)| \leq c \langle x \rangle^{2\alpha-1}.$$

Using the formula (3.8) with $\psi_N(x)$ instead of $\phi_N(x)$, and estimates similar to that in (3.9) it follows that

$$\lim_{m \uparrow \infty} \int_0^{t_1} \int (\partial_x u_m)^2 \psi'_N(x) dx dt = \int_0^{t_1} \int (\partial_x u)^2 \psi'_N(x) dx dt,$$

and

$$\left| \int_0^{t_1} \int (\partial_x u)^2 \psi'_N(x) dx dt \right| \leq \int_0^{t_1} \int (\partial_x u)^2 \langle x \rangle^{2\alpha-1} dx dt \leq M.$$

From these estimates and integrating in the time interval $[0, t] \subset [0, t_1]$ one obtains that

$$(3.18) \quad \langle x \rangle^\alpha u(t) \in L^2(\mathbb{R}) \quad t \in [0, t_1].$$

Combining (3.15) and (3.18) one sees that there exists $t^* \in [0, t_1]$ such that

$$\langle x \rangle^\alpha u(t^*), \partial_x u(t^*) \langle x \rangle^{\alpha-1/2} \in L^2(\mathbb{R}).$$

Hence,

$$(3.19) \quad \langle x \rangle^\alpha u(t^*), J(\langle x \rangle^{\alpha-1/2} u(\cdot, t^*)) = (1 - \partial_x^2)^{1/2} (\langle x \rangle^{\alpha-1/2} u(t^*)) \in L^2(\mathbb{R}).$$

Next we shall use the following lemma (see [1] and [12]) :

Lemma 3.1. *Let $a, b > 0$. Assume that $J^a f = (1 - \partial_x^2)^{a/2} f \in L^2(\mathbb{R})$ and $\langle x \rangle^b f = (1 + |x|^2)^{b/2} f \in L^2(\mathbb{R})$. Then for any $\theta \in (0, 1)$*

$$(3.20) \quad \|J^{\theta a} (\langle x \rangle^{(1-\theta)b} f)\|_2 \leq c \|\langle x \rangle^b f\|_2^{1-\theta} \|J^a f\|_2^\theta.$$

Defining

$$f = \langle x \rangle^{\alpha-1/2} u(t^*)$$

it follows that

$$Jf, \langle x \rangle^{1/2} f \in L^2(\mathbb{R})$$

and one gets from (3.20) with $\theta = 2\alpha$ that

$$(3.21) \quad J^{2\alpha} (\langle x \rangle^{1/2-\alpha} f) = J^{2\alpha} u(\cdot, t^*) \in L^2(\mathbb{R}).$$

Once (3.21) has been established the rest of the proof of Theorem 1.4 follows the argument described in [13].

Next, we consider the case $\alpha \in (1/2, 1]$.

From the previous case we already know that

$$u \in C([0, t_1] : H^1(\mathbb{R}) \cap L^2(\langle x \rangle dx)).$$

We also observe that for $\alpha \in (1/2, 1]$

$$\phi'_{N,\alpha}(x) + |\psi'_{N,\alpha}(x)| + |\phi_{N,\alpha-1/2}(x)| \leq c \langle x \rangle, \quad |\phi_{N,\alpha}^{(3)}(x)| \leq c.$$

As before we get

$$(3.22) \quad \frac{d}{dt} \int (u_m)^2 \phi_{N,\alpha} dx + 3 \int (\partial_x u_m)^2 \phi'_{N,\alpha} dx = \int (u_m)^2 \phi_{N,\alpha}^{(3)} dx + \frac{2}{3} \int (u_m)^3 \phi'_{N,\alpha} dx,$$

with $u_m \in C(\mathbb{R} : \mathcal{S}(\mathbb{R}))$ which justifies the integration by parts in (3.22).

In (3.22) we first use that

$$(3.23) \quad \left| \int (u_m)^2 \phi_{N,\alpha}^{(3)} dx \right| \leq c \|u_{0,m}\|_2^2 \leq 2c \|u_0\|_2^2.$$

Next, we estimate the cubic term in (3.22) (last term there). For this we observe that

$$(3.24) \quad |\phi'_{N,\alpha}(x)| \leq c |\phi_{N,\alpha-1/2}(x)|^2,$$

with c independent of N . Thus,

$$(3.25) \quad \left| \int (u_m)^3 \phi'_{N,\alpha}(x) dx \right| \leq c \|u_m(t)\|_\infty \|u_m(t) \phi_{N,\alpha-1/2}\|_2^2.$$

Combining the facts that for each N fixed the $\phi_{N,\alpha}$'s are bounded and that

$$\sup_{[0,t_1]} \|(u_m - u)(t)\|_2 \rightarrow 0 \quad \text{as } m \uparrow \infty,$$

one has

$$\sup_{[0,t_1]} \|(u_m - u)(t) \phi_{N,\alpha-1/2}\|_2 \rightarrow 0 \quad \text{as } m \uparrow \infty,$$

and consequently,

$$(3.26) \quad \sup_{[0,t_1]} \|u_m(t) \phi_{N,\alpha-1/2}\|_2 \leq 2 \sup_{[0,t_1]} \|u(t) \phi_{N,\alpha-1/2}\|_2 \leq 2 \sup_{[0,t_1]} \|\langle x \rangle^{1/2} u(t)\|_2 \leq M,$$

with $M = M(\|\langle x \rangle^{1/2} u_0\|_2, \|\langle x \rangle^{1/2} u(t_1)\|_2)$ for $m \gg 1$.

Inserting the above estimates in (3.22) and following the argument in the previous case $\alpha \in (0, 1/2]$ one gets that

$$(3.27) \quad \int_0^{t_1} \int (\partial_x u_m)^2(x, t) \phi'_N(x) dx dt \leq (1 + t_1)M,$$

for $m \gg 1$ and

$$(3.28) \quad \int_0^{t_1} \int (\partial_x u)^2(x, t) \langle x \rangle^{2\alpha-1} dx dt \leq (1 + t_1)M,$$

where $M = M(\|\langle x \rangle^{1/2} u_0\|_2, \|\langle x \rangle^{1/2} u(t_1)\|_2)$.

Next, we deduce (3.22) with $\psi_{N,\alpha}$ instead of $\varphi_{N,\alpha}$ and use (3.27) and (3.28) as in the previous case to get that

$$(3.29) \quad \langle x \rangle^\alpha u(t) \in L^2(\mathbb{R}) \quad t \in [0, t_1].$$

Once that the estimates (3.28) and (3.29) have been established the argument in the previous case involving Lemma 3.1 provides the desired result.

A similar boot-strap argument can be used in the case $\alpha \in (1, 3/2]$ and respectively for higher α .

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REFERENCES

- [1] J. Bergh and J. Löfström, *Interpolation spaces*, Springer New York and Berlin (1970).
- [2] J. Bourgain, *Fourier transform restriction phenomena for certain lattice subsets and applications to nonlinear evolution equations I, II*, *Geom. Funct. Anal.* **3** (1993) 107–156, 209–262.
- [3] J. Colliander, M. Keel, G. Staffilani, H. Takaoka, and T. Tao, *Sharp global wellposedness results for periodic and non-periodic KdV and modified KdV on \mathbb{R} and \mathbb{T}* , *J. Amer. Math. Soc.* **16** (2003), 705–749.
- [4] L. Escauriaza, C. E. Kenig, G. Ponce, and L. Vega, *On Uniqueness Properties of Solutions of the k -generalized KdV*, *J. Funct. Anal.* **244**, 2 (2007), 504–535.
- [5] J. Ginibre and Y. Tsutsumi, *Uniqueness for the generalized Korteweg-de Vries equations*, *SIAM J. Math. Anal.* **20** (1989), 1388–1425.
- [6] A. Grünrock, *A bilinear Airy type estimate with application to the 3-gkdv equation*, *Diff. Int. Eqs.* **18** (2005), 1333–1339.
- [7] A. Grünrock, M. Panthee, and J. Drumond Silva, *A remark on global well-posedness below L^2 for the gKdV-3 equation*, *Diff. Int. Eqs.* **20** (2007), 1229–1236.
- [8] T. Kato, *On the Cauchy problem for the (generalized) Korteweg-de Vries equation*, *Advances in Mathematics Supplementary Studies, Studies in Applied Math.* **8** (1983), 93–128.
- [9] C. E. Kenig, G. Ponce, and L. Vega, *Well-posedness and scattering results for the generalized Korteweg-de Vries equation via contraction principle*, *Comm. Pure Appl. Math.* **46** (1993), 527–620.
- [10] C. E. Kenig, G. Ponce, and L. Vega, *A bilinear estimate with applications to the KdV equation*, *J. Amer. Math. Soc.* **9** (1996), 573–603.
- [11] S. N. Kruzhkov and A. V. Faminskii, *Generalized solutions of the Cauchy problem for the Korteweg-de Vries equation*, *Math. U.S.S.R. Sbornik* **48** (1984), 93–138.
- [12] J. Nahas and G. Ponce, *On the persistent properties of solutions to semi-linear Schrödinger equation*, *Comm. P.D.E.* **34** (2009), 1–20.
- [13] J. Nahas and G. Ponce, *On the well-posedness of the modified Korteweg-de Vries equation in weighted Sobolev spaces*, preprint.
- [14] L. Robbiano, *Unicité forte à l’infini pour KdV*, *Control Opt. and Cal. Var.* **8** (2002), 933–939.
- [15] T. Tao, *Scattering for the quartic generalized Korteweg-de Vries equation*, *J. Differential Equations* **232** (2007), 623–651.

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