

DRM 9/6/18  
 supersymmetry &  
 Seiberg-Witten Theory

4D,  $\mathcal{N}=2$   $SU(2)$  gauge theory

(no matter) following

Seiberg-Witten hep-th/9407087

$\mathcal{N}=2$  multiplets:

$\mathcal{N}=2$  vector: gauge field  $A_\mu$   
 fermions in adjoint rep of  $G$ : two Weyl fermions  $\lambda, \bar{\lambda}$   
 - scalar  $\phi$

$G = SU(2)$

$\mathcal{N}=2$  hypermultiplet:

fermions in 2 reps of  $G$ : two Weyl fermions  $\psi, \bar{\psi}$   
 two scalars  $\varphi, \tilde{\varphi}$

In  $\mathcal{N}=1$  terms

$\mathcal{N}=2$  vector = ( $\mathcal{N}=1$ ) vector + ( $\mathcal{N}=1$ ) chiral

$\mathcal{N}=2$  hyper = 2 ( $\mathcal{N}=1$ ) chirals

Classical R symmetry:

different actions on fermions:  $SU(2)_R$  acts on  $\{\psi, \bar{\psi}\}$   
 $U(1)_R$  acts on  $\phi$

This  $U(1)_R$  is (almost always) anomalous.

$G = SU(N_c)$  and  $N_f$  "fundamental" hypermultiplets

$\Rightarrow$  anomaly free subgroup is  $Z_{4N_c - 2N_f} \subseteq U(1)_R$

Classical potential for  $\mathcal{N}=2$  theory:

( $g$  coupling constant)

$V(\phi) = \frac{1}{g^2} \text{Tr}[\phi, \phi^\dagger]^2$

$V(\phi) = 0$  can happen if

$[\phi, \phi^\dagger] = 0$

$\rightarrow$  space of solutions (vacuum degeneracy)

( $\phi$  takes values in the algebra)

$G = SU(2)$ ,  $\phi$  is conjugate to

$\frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

$u = \frac{1}{2} a^2 = \text{Tr} \phi^2$

is the natural parameter.

$a \neq 0 \Rightarrow SU(2)$  breaks to commutant of

$\frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$  which is  $U(1)$

(weyl  $SU(2)$   $a \mapsto -a$ )



promote the parameter to a field

$A = (W=2)$  vect mult, whose  
scalar comp is a

Terms in susy low-energy  
theory involve a holo function

$$F(A).$$

In  $N=1$  language  
the lagrangian:

$$\frac{1}{4\pi} \text{Im} \left[ \int d^4\theta \frac{\partial F(A)}{\partial A} A \right. \\ \left. + \int d^2\theta \frac{1}{2} \frac{\partial^2 F(A)}{\partial A^2} W_\alpha W^\alpha \right]$$

Structure of the theory is encoded  
in  $F(A)$

$$\tau_{cl} = \frac{e}{2\pi} + i \frac{4\pi}{g^2}$$

Kahler potential for matrix in space of  
scalars:

$$K = \text{Im} \left( \frac{\partial F}{\partial A} \bar{A} \right)$$

$$ds^2 = \text{Im} \frac{\partial^2 F(a)}{\partial a^2} da d\bar{a}$$

= kahler metric on the  
space of vacua.

Features of the quantum theory:

- 1) If  $|c| \gg 0$ , theory is  
"asymptotically free"  
(weakly coupled)
- 2) vacuum degeneracy cannot be  
removed quantum mechanically  
since  $N=2$  supersymmetry does  
not permit an  $N=1$  superpotential
- 3) classical value of  $F$  is

$$F(A) = \frac{1}{2} \tau_{cl} A^2$$

$$\tau(a) = \frac{\partial^2 F}{\partial a^2}$$

$$F_{one-loop} = i \frac{1}{2\pi} A^2 \ln \frac{A^2}{\Lambda^2}$$

( $\Lambda$  = dimensional  
scale)

( $N=2$  susy  $\Rightarrow$  1 loop exact)

$$F = i \frac{1}{2\pi} A^2 \ln \frac{A^2}{\Lambda^2}$$

$$+ \sum_{k=2}^{\infty} F_k \left( \frac{A}{\Lambda} \right)^k A^2$$

(holomorphic corrections)

$$F_1 \neq 0$$

Duality:  $ds^2 = \text{Im}(\tau) da d\bar{a}$

for large  $|\tau|$

$\tau(\sigma) \approx i \left( \ln\left(\frac{\sigma^2}{\Lambda^2}\right) + 3 \right)$

$\tau(\sigma)$  multi-valued

$ds^2$  is not pos. definite  $\forall a$ .

let  $a_D = \frac{\partial F}{\partial a}$

$ds^2 = \text{Im}(da_D d\bar{a})$

$= -\frac{i}{2} (da_n d\bar{a} - da d\bar{a}_D)$

Observation:

$a \leftrightarrow -a_D$  gives another description of the same theory.

Note:  $a_D \rightarrow a_D + b_C$  leaves the invariant  
 $a \rightarrow a$

$\left( \begin{matrix} 1 & b \\ 0 & 1 \end{matrix} \right) \left( \begin{matrix} a_D \\ a \end{matrix} \right), \left( \begin{matrix} 0 & 1 \\ -1 & 0 \end{matrix} \right) \left( \begin{matrix} a_D \\ a \end{matrix} \right)$

$SL(2, \mathbb{R})$   
 $U(1)$  broken to  
 $SL(2, \mathbb{Z})$  by BPS particles

$z = a\eta_L + c_0 \eta_M$

light magnetic  $e_D = 0$



(require at least 4 singularities to give multi-valuedness)

light electric

$a \neq 0$

$a \rightarrow -a$

action of Weyl ( $SU(2)$ )

$u = \frac{1}{2} a^2$