

Random Blob Methods for Diffusion

2026 SIAM Annual Meeting

Claire Murphy

UC Santa Barbara, joint work with Dr. Katy Craig

July 10, 2026



The Nonlinear Fokker-Planck Equation

Suppose that

- $v : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}^d$,
- $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is l.s.c., coercive, convex, and $f(0) = 0$.

Consider the **nonlinear Fokker-Planck equation**

$$\begin{cases} \partial_t \rho = \nabla \cdot (\nabla f^*(\rho) + \rho v) \\ \rho \in \partial f(\rho) \\ \rho(0) = \rho_0. \end{cases} \quad (\text{PDE})$$

The Nonlinear Fokker-Planck Equation

Suppose that

- $v : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}^d$,
- $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is l.s.c., coercive, convex, and $f(0) = 0$.

Consider the **nonlinear Fokker-Planck equation**

$$\begin{cases} \partial_t \rho = \nabla \cdot (\nabla f^*(\rho) + \rho v) \\ \rho \in \partial f(\rho) \\ \rho(0) = \rho_0. \end{cases} \quad (\text{PDE})$$

This equation describes cases of linear and nonlinear diffusion, such as:

- 0 The Fokker-Planck equation: $f(s) = s \ln(s) - s \Rightarrow \partial_t \rho = \Delta \rho + \nabla \cdot (\rho v)$.

The Nonlinear Fokker-Planck Equation

Suppose that

- $v : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}^d$,
- $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is l.s.c., coercive, convex, and $f(0) = 0$.

Consider the **nonlinear Fokker-Planck equation**

$$\begin{cases} \partial_t \rho = \nabla \cdot (\nabla f^*(\rho) + \rho v) \\ \rho \in \partial f(\rho) \\ \rho(0) = \rho_0. \end{cases} \quad (\text{PDE})$$

This equation describes cases of linear and nonlinear diffusion, such as:

- 0 The Fokker-Planck equation: $f(s) = s \ln(s) - s \Rightarrow \partial_t \rho = \Delta \rho + \nabla \cdot (\rho v)$.
- 1 The porous media / fast diffusion equation [Ott, Vaz]: If $m > 1 - \frac{2}{d+2}$,

$$f(s) = \begin{cases} s \ln(s) - s & m = 1 \\ \frac{s^m}{m-1} & m \neq 1 \end{cases} \Rightarrow \partial_t \rho = \Delta \rho^m + \nabla \cdot (\rho v).$$

The Nonlinear Fokker-Planck Equation

Suppose that

- $v : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}^d$,
- $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is l.s.c., coercive, convex, and $f(0) = 0$.

Consider the **nonlinear Fokker-Planck equation**

$$\begin{cases} \partial_t \rho = \nabla \cdot (\nabla f^*(\rho) + \rho v) \\ \rho \in \partial f(\rho) \\ \rho(0) = \rho_0. \end{cases} \quad (\text{PDE})$$

This equation describes cases of linear and nonlinear diffusion, such as:

- 0 The Fokker-Planck equation: $f(s) = s \ln(s) - s \Rightarrow \partial_t \rho = \Delta \rho + \nabla \cdot (\rho v)$.
- 1 The porous media / fast diffusion equation [Ott, Vaz]: If $m > 1 - \frac{2}{d+2}$,

$$f(s) = \begin{cases} s \ln(s) - s & m = 1 \\ \frac{s^m}{m-1} & m \neq 1 \end{cases} \Rightarrow \partial_t \rho = \Delta \rho^m + \nabla \cdot (\rho v).$$

- 2 Height-constrained transport [DMSV, AKY]:

$$f(s) = \begin{cases} 0 & 0 \leq s \leq 1 \\ \infty & \text{otherwise} \end{cases} \Rightarrow \begin{cases} \partial_t \rho + \nabla \cdot (\rho v) = 0 \\ 0 \leq \rho \leq 1. \end{cases} \quad "$$

Longtime Behavior

Setup: Let $\tilde{\rho}$ be a probability measure on Euclidean space $\mathbb{R}^d$. Assume $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is differentiable and $v = \nabla V$, where $V : \mathbb{R}^d \rightarrow \mathbb{R}$ is continuously differentiable and λ -convex for $\lambda > 0$.

Longtime Behavior

Setup: Let $\tilde{\rho}$ be a probability measure on Euclidean space $\mathbb{R}^d$. Assume $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is differentiable and $v = \nabla V$, where $V : \mathbb{R}^d \rightarrow \mathbb{R}$ is continuously differentiable and λ -convex for $\lambda > 0$.

1 Linear Fokker-Planck equation

1 **Assumption:** $\tilde{\rho}(x) = Ze^{-V(x)} dx$. Z is a constant chosen s.t. $\int \tilde{\rho} = 1$.

2 **PDE perspective:** [AGS] If $\rho(t)$ is a solution to the **Fokker-Planck equation** with $v = \nabla V$, then

$$W_2(\rho(t), \tilde{\rho}) \leq e^{-\lambda t} W_2(\rho_0, \tilde{\rho}).$$

Longtime Behavior

Setup: Let $\tilde{\rho}$ be a probability measure on Euclidean space $\mathbb{R}^d$. Assume $f : [0, \infty) \rightarrow \mathbb{R} \cup \infty$ is differentiable and $v = \nabla V$, where $V : \mathbb{R}^d \rightarrow \mathbb{R}$ is continuously differentiable and λ -convex for $\lambda > 0$.

1 Linear Fokker-Planck equation

1 Assumption: $\tilde{\rho}(x) = Ze^{-V(x)} dx$. Z is a constant chosen s.t. $\int \tilde{\rho} = 1$.

2 PDE perspective: [AGS] If $\rho(t)$ is a solution to the **Fokker-Planck equation** with $v = \nabla V$, then

$$W_2(\rho(t), \tilde{\rho}) \leq e^{-\lambda t} W_2(\rho_0, \tilde{\rho}).$$

2 Nonlinear Fokker-Planck equation

1 Assumption: $\tilde{\rho}(x) = \max((f')^{-1}(Z - V(x)), 0)$. Z is a constant chosen s.t. $\int \tilde{\rho} = 1$.

2 PDE perspective: [CJT] Formally, if $\rho(t)$ is a solution to the **nonlinear Fokker-Planck equation** with $v = \nabla V$, then

$$W_2(\rho(t), \tilde{\rho}) \leq e^{-\lambda t} W_2(\rho_0, \tilde{\rho}).$$

Blob Methods for the Nonlinear Fokker-Planck Equation

In recent years, blob methods have shown tremendous promise in simulating a wide range of PDEs [BE, CCP, CCY, CESW, CEHT, CJT]. The main challenge in the implementation of these methods is their **computational complexity**.

Blob Methods for the Nonlinear Fokker-Planck Equation

In recent years, blob methods have shown tremendous promise in simulating a wide range of PDEs [BE, CCP, CCY, CESW, CEHT, CJT]. The main challenge in the implementation of these methods is their **computational complexity**.

[CJT] showed convergence of a variant of the **regularized** PDE

$$\begin{cases} \partial_t \rho - \nabla \cdot (\rho [f_\varepsilon''(\varphi_\varepsilon * \rho)(\nabla \varphi_\varepsilon * \rho) + v]) = 0 \\ \rho(0, x) = \rho_0(x). \end{cases} \quad (\text{PDE}_\varepsilon)$$

to (PDE) as $\varepsilon \rightarrow 0$. Here, φ_ε is a mollifier and f_ε is a suitable regularization of f .

Blob Methods for the Nonlinear Fokker-Planck Equation

In recent years, blob methods have shown tremendous promise in simulating a wide range of PDEs [BE, CCP, CCY, CESW, CEHT, CJT]. The main challenge in the implementation of these methods is their **computational complexity**.

[CJT] showed convergence of a variant of the **regularized** PDE

$$\begin{cases} \partial_t \rho - \nabla \cdot (\rho [f''_\varepsilon(\varphi_\varepsilon * \rho)(\nabla \varphi_\varepsilon * \rho) + v]) = 0 \\ \rho(0, x) = \rho_0(x). \end{cases} \quad (\text{PDE}_\varepsilon)$$

to (PDE) as $\varepsilon \rightarrow 0$. Here, φ_ε is a mollifier and f_ε is a suitable regularization of f .

If $\rho_0(x) = \sum_{i=1}^N m_i \delta_{x_i^0}$, then $\rho(t, x) = \sum_{i=1}^N m_i \delta_{x_i(t)}$, where each x_i moves according to the ODE system

$$\begin{cases} \dot{x}_i(t) = -f''_\varepsilon \left(\sum_{j=1}^N m_j \varphi_\varepsilon(x_i(t) - x_j(t)) \right) \sum_{j=1}^N m_j \nabla \varphi_\varepsilon(x_i(t) - x_j(t)) - v(x_i(t)) \\ x_i(0) = x_i^0. \end{cases} \quad (1)$$

Blob Methods for the Nonlinear Fokker-Planck Equation

In recent years, blob methods have shown tremendous promise in simulating a wide range of PDEs [BE, CCP, CCY, CESW, CEHT, CJT]. The main challenge in the implementation of these methods is their **computational complexity**.

[CJT] showed convergence of a variant of the **regularized** PDE

$$\begin{cases} \partial_t \rho - \nabla \cdot (\rho [f''_\varepsilon (\varphi_\varepsilon * \rho) (\nabla \varphi_\varepsilon * \rho) + v]) = 0 \\ \rho(0, x) = \rho_0(x). \end{cases} \quad (\text{PDE}_\varepsilon)$$

to (PDE) as $\varepsilon \rightarrow 0$. Here, φ_ε is a mollifier and f_ε is a suitable regularization of f .

If $\rho_0(x) = \sum_{i=1}^N m_i \delta_{x_i^0}$, then $\rho(t, x) = \sum_{i=1}^N m_i \delta_{x_i(t)}$, where each x_i moves according to the ODE system

$$\begin{cases} \dot{x}_i(t) = -f''_\varepsilon \left(\sum_{j=1}^N m_j \varphi_\varepsilon(x_i(t) - x_j(t)) \right) \sum_{j=1}^N m_j \nabla \varphi_\varepsilon(x_i(t) - x_j(t)) - v(x_i(t)) \\ x_i(0) = x_i^0. \end{cases} \quad (1)$$

Our goal is to improve the computational efficiency of blob methods.

Previous Work: The Random Batch Method

The RB method has been successfully used to solve ODE systems that model

- 1 Opinion dynamics [PZ],
- 2 The movement of plasma particles [JXL],
- 3 Statistical sampling and machine learning [JL].

Previous Work: The Random Batch Method

The RB method has been successfully used to solve ODE systems that model

- 1 Opinion dynamics [PZ],
- 2 The movement of plasma particles [JXL],
- 3 Statistical sampling and machine learning [JL].

At each time step, particles are split into random “batches” of size p . If I_i equals the index set of all particles in the same batch as particle i ,

$$x_i^{m+1} = x_i^m + \Delta t \left(-f''_\varepsilon \left(\sum_{j \in I_i} \tilde{m}_j \varphi_\varepsilon(x_i^m - x_j^m) \right) \sum_{j \in I_i} \tilde{m}_j \nabla \varphi_\varepsilon(x_i^m - x_j^m) - v(x_i^m) \right),$$

where $\tilde{m}_j = m_j / \sum_{k \in I_i} m_k$ [JLL].

Previous Work: The Random Batch Method

The RB method has been successfully used to solve ODE systems that model

- 1 Opinion dynamics [PZ],
- 2 The movement of plasma particles [JXL],
- 3 Statistical sampling and machine learning [JL].

At each time step, particles are split into random “batches” of size p . If I_i equals the index set of all particles in the same batch as particle i ,

$$x_i^{m+1} = x_i^m + \Delta t \left(-f''_\varepsilon \left(\sum_{j \in I_i} \tilde{m}_j \varphi_\varepsilon(x_i^m - x_j^m) \right) \sum_{j \in I_i} \tilde{m}_j \nabla \varphi_\varepsilon(x_i^m - x_j^m) - v(x_i^m) \right),$$

where $\tilde{m}_j = m_j / \sum_{k \in I_i} m_k$ [JLL].

	Runtime	Error for fixed N
FE	$O(N^2/\Delta t)$	$O(\Delta t)$
RB	$O(Np/\Delta t)$	$O(\sqrt{\Delta t/p})$

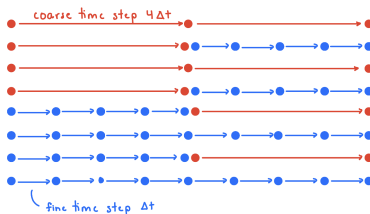
Table: Δt refers to the time step size, N refers to the number of particles, and p refers to the batch size.

As we will see, the random batch method is not well-suited to the ODE system (1).

Our Contribution: The Random Multirate Method

The Random Multirate Method:

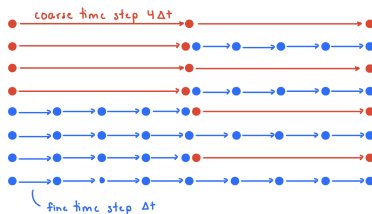
- 1 $N - p$ randomly selected particles take a FE step with coarse step size $k\Delta t$.
- 2 Remaining p particles take k FE steps with fine step size Δt .



Our Contribution: The Random Multirate Method

The Random Multirate Method:

- 1 $N - p$ randomly selected particles take a FE step with coarse step size $k\Delta t$.
- 2 Remaining p particles take k FE steps with fine step size Δt .



	Runtime	Error for fixed N
FE	$O(N^2/\Delta t)$	$O(\Delta t)$
RB	$O(Np/\Delta t)$	$O(\sqrt{\Delta t/p})$
RM	$O\left(\frac{N}{\Delta t} \left(p + \frac{N-p}{k}\right)\right)$	$O(k\Delta t)$

Table: Δt refers to the time step size, N refers to the number of particles, p refers to the batch size, and k refers to the time step ratio.

Height Constraint Results

Height constraint: Our continuum solution cannot rise above height 1.

$$f(s) = \begin{cases} 0 & s \in [0, 1] \\ \infty & \text{otherwise} \end{cases} \Rightarrow \begin{cases} \partial_t \rho + \nabla \cdot (\rho v) = 0 \\ 0 \leq \rho \leq 1. \end{cases}$$

Note that $f_m(s) := \frac{s^m}{m-1}$ converges to f as $m \rightarrow \infty$. We approximate f with f_m , $m = 100$.

Height Constraint Results

Height constraint: Our continuum solution cannot rise above height 1.

$$f(s) = \begin{cases} 0 & s \in [0, 1] \\ \infty & \text{otherwise} \end{cases} \Rightarrow \begin{cases} \partial_t \rho + \nabla \cdot (\rho v) = 0 \\ 0 \leq \rho \leq 1. \end{cases}$$

Note that $f_m(s) := \frac{s^m}{m-1}$ converges to f as $m \rightarrow \infty$. We approximate f with f_m , $m = 100$.

Continuum behavior with external velocity $v(x) = x$.

- 1 For $t \leq \ln(3) \approx 1.098$, the solution is attracted to $x_0 = 0$.
- 2 After the solution hits the "ceiling" 1, it remains steady.

Height Constraint Results

Height constraint: Our continuum solution cannot rise above height 1.

$$f(s) = \begin{cases} 0 & s \in [0, 1] \\ \infty & \text{otherwise} \end{cases} \Rightarrow \begin{cases} \partial_t \rho + \nabla \cdot (\rho v) = 0 \\ 0 \leq \rho \leq 1. \end{cases}$$

Note that $f_m(s) := \frac{s^m}{m-1}$ converges to f as $m \rightarrow \infty$. We approximate f with f_m , $m = 100$.

Continuum behavior with external velocity $v(x) = x$.

- 1 For $t \leq \ln(3) \approx 1.098$, the solution is attracted to $x_0 = 0$.
- 2 After the solution hits the “ceiling” 1, it remains steady.

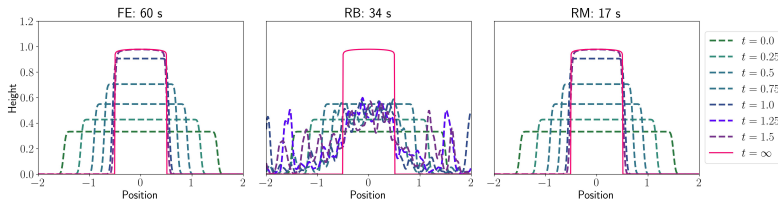


Figure: $N = 600$, $\Delta t = 10^{-3}$, $p = 300$, $k = 2$, $\varepsilon = .02108$ for φ_ε a Gaussian with standard deviation ε .

Height Constraint Continued

For different choice of Δt , how is the runtime and accuracy of the FE, RB, and RM methods affected?

We measure accuracy using an approximation of the 2-Wasserstein distance between the particle solution and the exact continuum solution at final time $T = 1.5$.

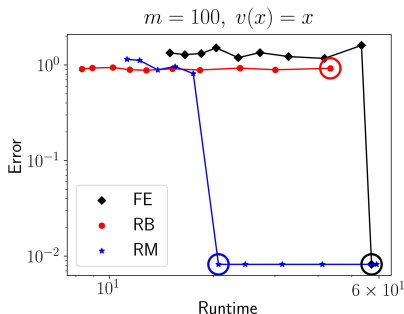


Figure: $N = 600, p = 300, k = 2, \varepsilon = 0.02108$ for φ_ε a Gaussian with standard deviation ε .

Diffusion Regime Results

Diffusion: We simulate nonlinear diffusion with

$$f(s) = \begin{cases} s \ln(s) - s & m = 1 \\ \frac{s^m}{m-1} & m \neq 1 \end{cases} \Rightarrow \partial_t \rho = \Delta \rho^m + \nabla \cdot (\rho v)$$

and $v(x) = \frac{x}{m+1}$ up to final time $T = 1$.

Diffusion Regime Results

Diffusion: We simulate nonlinear diffusion with

$$f(s) = \begin{cases} s \ln(s) - s & m = 1 \\ \frac{s^m}{m-1} & m \neq 1 \end{cases} \Rightarrow \partial_t \rho = \Delta \rho^m + \nabla \cdot (\rho v)$$

and $v(x) = \frac{x}{m+1}$ up to final time $T = 1$.

Again, we measure accuracy using an approximation of the 2-Wasserstein distance between the particle solution and the exact continuum solution at final time $T = 1$.

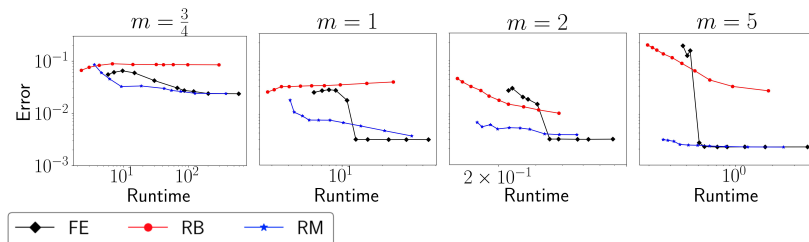


Figure: The performance of the FE, RB, and RM methods in the case $p = N/2$ and $k = 2$.

2D Results: The Navier-Stokes Equation

We use the RM method to simulate the viscosity formulation of the 2D Navier-Stokes equation

$$\partial_t \rho + v \cdot \nabla \rho = \nu \Delta \rho,$$

where $v = K * \rho$ for K the Biot-Savart kernel. [Nor] utilized regridding to simulate this equation with a “double bump” initial condition. We avoid regridding, instead using the RM method to simulate the solution.

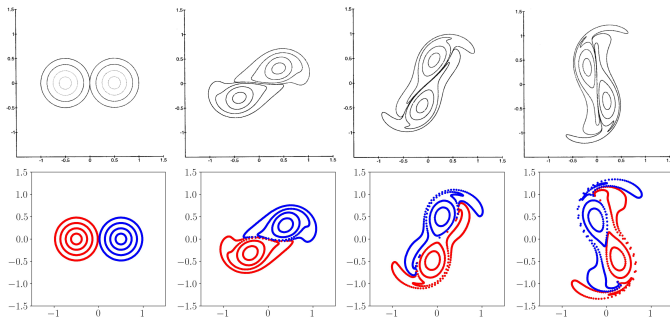


Figure: $N = 4688$, $\Delta t = 0.05$, $p = 2344$, $k = 2$, $\nu = 0.0005$ $\varepsilon = .08319$ for φ_ε a Gaussian with standard deviation ε .

Conclusions

Major takeaways

- 1 The RM method improves computational efficiency and preserves accuracy of blob methods for diffusion equations.
- 2 We see best performance of the RM method in the slow diffusion regime.
- 3 The RM method is compatible with a wide choice of linear and nonlinear diffusion equations.

Conclusions

Major takeaways

- 1 The RM method improves computational efficiency and preserves accuracy of blob methods for diffusion equations.
- 2 We see best performance of the RM method in the slow diffusion regime.
- 3 The RM method is compatible with a wide choice of linear and nonlinear diffusion equations.

Directions for future work

- 1 Mean field variational inference
- 2 Consensus-based optimization

Conclusions

Major takeaways

- 1 The RM method improves computational efficiency and preserves accuracy of blob methods for diffusion equations.
- 2 We see best performance of the RM method in the slow diffusion regime.
- 3 The RM method is compatible with a wide choice of linear and nonlinear diffusion equations.

Directions for future work

- 1 Mean field variational inference
- 2 Consensus-based optimization

THANK YOU!!!

References

- AGS Ambrosio, Luigi, Nicola Gigli, and Giuseppe Savaré. "Gradient flows in metric spaces and in the space of probability measures." Birkhäuser, 2008.
- AKY Alexander, Damon, Inwon Kim, and Yao Yao. "Quasi-static evolution and congested crowd transport." Nonlinearity, 2014.
- BE Burger, Martin and Antonio Esposito. "Porous medium equation and cross-diffusion systems as limit of nonlocal interaction." Nonlinear Analysis, 2022.
- BJ Bántay, Péter and Imre M. Jánosi. "Avalanche dynamics from anomalous diffusion." Physical Review Letters, 1992.
- CCP Carrillo, José Antonio, Katy Craig, and Francesco Patacchini. "A blob method for diffusion." Calculus of Variations and PDEs, 2018.
- CCY Carrillo, José Antonio, Katy Craig, and Yao Yao. "Aggregation-diffusion equations: Dynamics, asymptotics, and singular limits." Active Particles, volume 2, 2019.

References cont'd

- CESW** Carrillo, José Antonio, et al. "Nonlocal particle approximations for linear and fast diffusion equations." Preprint, 2024.
- CEHT** Craig, Katy et al. "A blob method for inhomogeneous diffusion with applications to multi-agent control and sampling." *Mathematics of Computation*, 2023.
- CJT** Craig, Katy, Matt Jacobs, and Olga Turanova. "Nonlocal approximation of slow and fast diffusion." *Journal of Differential Equations*, 2025.
- DMSV** de Philippis, Guido et al. "BV Estimates in Optimal Transportation and Applications." *Archive for Rational Mechanics and Analysis*, 2016.
- JL** Jin, Shi and Lei Li. "Random Batch Methods for Classical and Quantum Interacting Particle Systems and Statistical Samplings." *Modeling and Simulation in Science, Engineering and Technology*, 2021.
- JLL** Jin, Shi, Lei Li, and Jian-Guo Liu. "Random batch methods (RBM) for interacting particle systems." *Journal of Computational Physics*, 2020.

References cont'd

- JXL** Jin, Shi and Xiantao Li. "Random Batch Algorithms for Quantum Monte Carlo Simulations." Communications in Computational Physics, 2020.
- Nor** Nordmark, Hans Olsson. "Deterministic high order vortex methods for the 2D Navier-Stokes equation with rezoning." Journal of Computational Physics, 1996.
- Ott** Otto, Felix. "The geometry of dissipative evolution equations: the porous medium equation." Communications in Partial Differential Equations, 2001.
- PZ** Pareschi, Lorenzo and Mattia Zanella. "Reduced Variance Random Batch Methods for Nonlocal PDEs." Acta Applicandae Mathematicae, 2024.
- Vaz** Vazquez, Juan Luis. "The porous medium equation." Oxford Mathematical Monographs, 2007.