

Celestial Holography from Bottom-up to Top-down

The Western Hemisphere Colloquium on Geometry and Physics

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Soft Thm=Ward Id

u-falloffs and antipodal matching

Celestial OPE

analytic structure of amplitudes, inner product

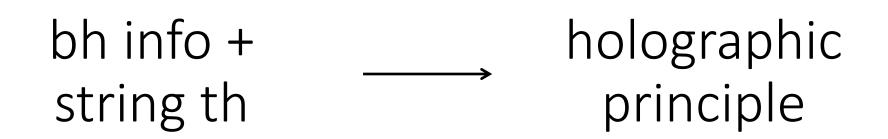
Boost Basis

principal series vs highest weight reps

Chiral Algebra

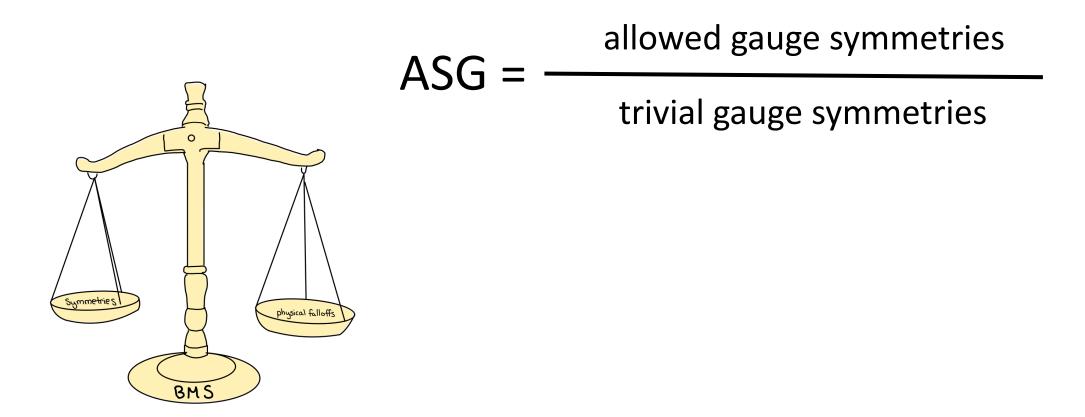
twistor & twisted holography re-interpretations

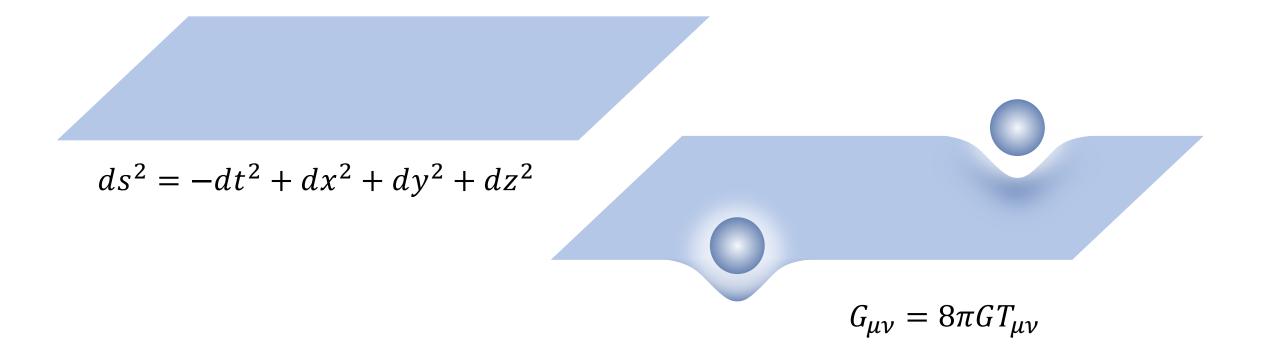
WHAT IS CELESTIAL HOLOGRAPHY?

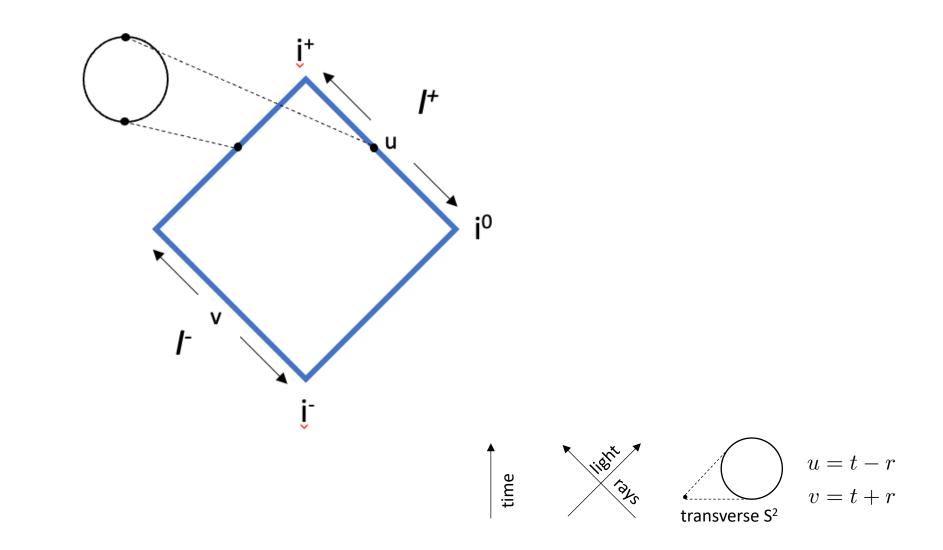


$\Lambda = 0 \qquad \qquad \text{vs} \qquad \Lambda \to 0$

Lesson 1: BMS >> Poincare: $\Lambda = 0$ spacetimes have a much larger class of possible symmetries.







In Bondi gauge the metric near future null infinity takes the form

$$ds^{2} = -du^{2} - 2dudr + 2r^{2}\gamma_{z\bar{z}}dzd\bar{z} + 2\frac{m_{B}}{r}du^{2} + (rC_{zz}dz^{2} + D^{z}C_{zz}dudz + \frac{1}{r}(\frac{4}{3}N_{z} - \frac{1}{4}\partial_{z}(C_{zz}C^{zz}))dudz + c.c.) + \dots$$
Radiative Data

which is preserved by the residual diffeomorphisms

$$\xi^{+} = (1 + \frac{u}{2r})Y^{+z}\partial_{z} - \frac{u}{2r}D^{\bar{z}}D_{z}Y^{+z}\partial_{\bar{z}} - \frac{1}{2}(u+r)D_{z}Y^{+z}\partial_{r} + \frac{u}{2}D_{z}Y^{+z}\partial_{u} + c.c$$

$$+ f^{+}\partial_{u} - \frac{1}{r}(D^{z}f^{+}\partial_{z} + D^{\bar{z}}f^{+}\partial_{\bar{z}}) + D^{z}D_{z}f^{+}\partial_{r}$$
Superrotations
Supertranslations

Form this bulk analysis we land on the BMS group

$$[L_m, L_n] = (m - n)L_{m+n}, \quad [\bar{L}_m, \bar{L}_n] = (m - n)\bar{L}_{m+n},$$
$$[L_n, P_{k,l}] = (\frac{1}{2}n - k)P_{k+n,l}, \quad [\bar{L}_n, P_{k,l}] = (\frac{1}{2}n - l)P_{k,l+n},$$
$$[P_{m,n}, P_{k,l}] = 0.$$

where $n, m \in \{-1, 0, 1\}$ & $k, l \in \left\{-\frac{1}{2}, \frac{1}{2}\right\}$ give the Poincare subalgebra

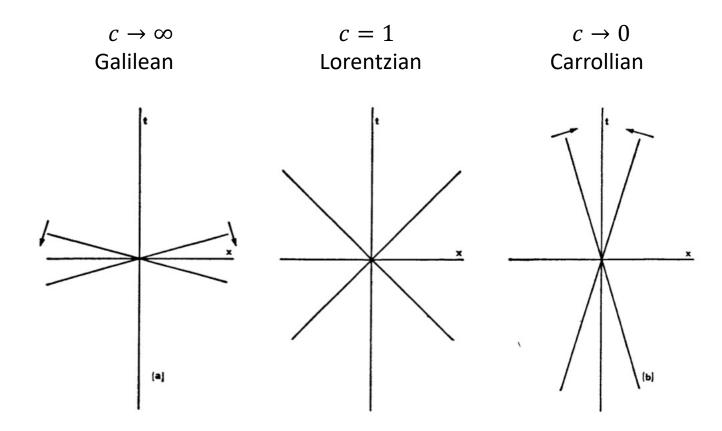
We can see this enhancement from the boundary perspective

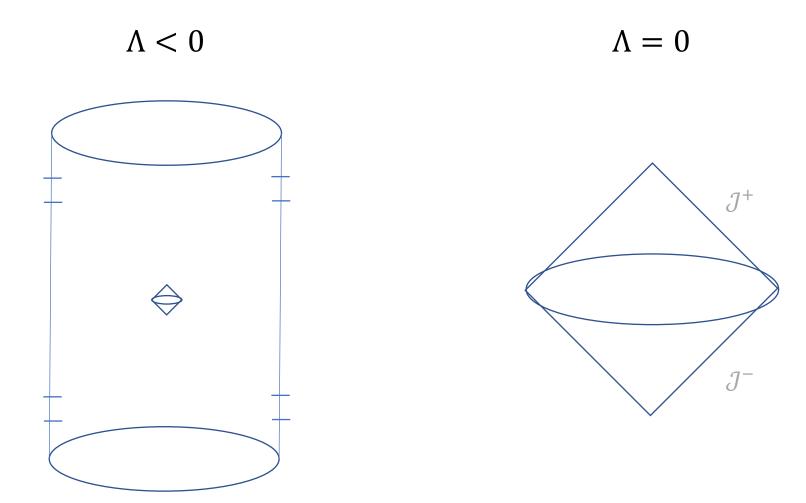
$$ds^2 = -c^2 du^2 + dz d\bar{z}$$

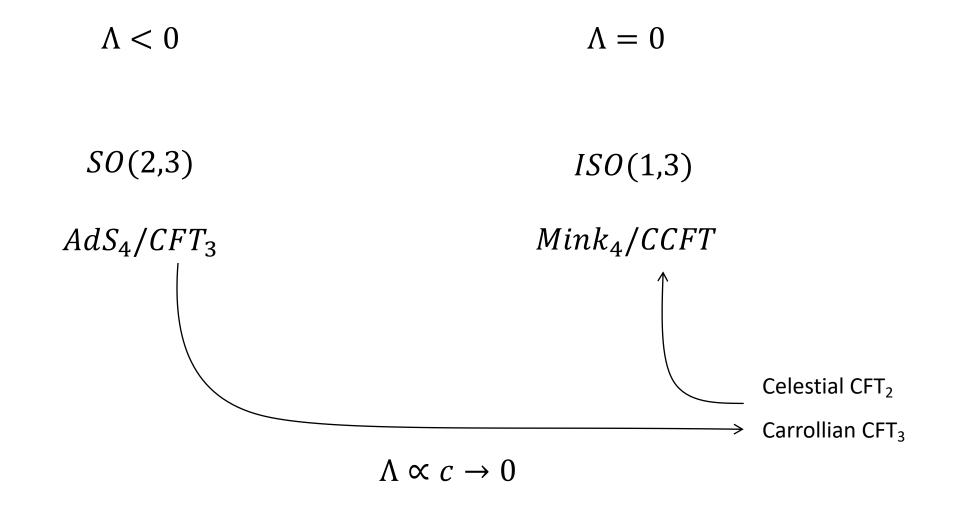
as a Carrollian limit of of a Lorentzian CFT₃.

$$\lim_{c \to 0} [\nabla_{\mu} \xi_{\nu} + \nabla_{\nu} \xi_{\mu} = \alpha g_{\mu\nu}]$$
$$\lim_{c \to 0} c^2 [\nabla^{\mu} \xi^{\nu} + \nabla^{\nu} \xi^{\mu} = -\alpha g^{\mu\nu}]$$

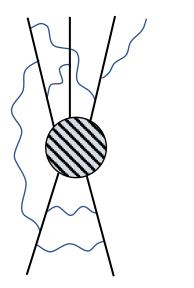
$$ds^2 = -c^2 dt^2 + d\vec{x}^2$$







Lesson 2: These are indeed symmetries of the perturbative S-matrix!



 \Leftrightarrow

Soft Thm = Ward Id

$$\langle out|Q^+[Y]\mathcal{S} - \mathcal{S}Q^-[Y]|in\rangle = 0$$

The free data for our solution takes the form

$$\{m_B(u_0, z, \bar{z}), N_z(u_0, z, \bar{z}), C_{zz}(u, z, \bar{z})\}$$

where the u dependence of the Bondi mass and angular momentum aspect are fixed by the constraint equations

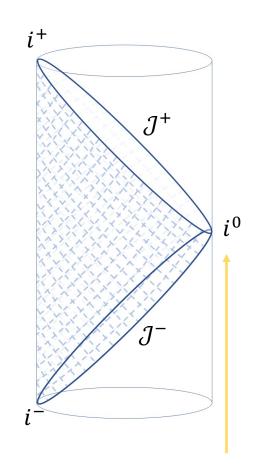
$$n^{\mu}[G_{\mu\nu} - 8\pi G T_{\mu\nu}] = 0$$

The free data for our solution takes the form

$$\{m_B(u_0, z, \bar{z}), N_z(u_0, z, \bar{z}), C_{zz}(u, z, \bar{z})\}$$

where the u dependence of the Bondi mass and angular momentum aspect are fixed by the constraint equations

$$\partial_u m_B = \frac{1}{4} [D_z^2 N^{zz} + D_{\bar{z}}^2 N^{\bar{z}\bar{z}}] - T_{uu}$$



1. Q $^+$ = Q $^-$ due to antipodal matching of m_B & N_z across i^0

2. Integration by parts turns the charges into fluxes

$$\begin{split} Q_{f}^{+} &= \frac{1}{8\pi G} \int_{\mathcal{F}_{-}^{+}} 2m_{B}f \\ & 8\pi G Q^{+}[Y] = \int_{\mathcal{I}^{+}} \sqrt{\gamma} d^{2}z du \ \left[-\frac{1}{2} D_{z}^{3} Y^{z} u \partial_{u} C^{zz} + Y^{z} T_{uz} + u D_{z} Y^{z} T_{uu} + h.c. \right] \\ & Q^{+}[Y] = Q_{S}^{+}[Y] + Q_{H}^{+}[Y] \end{split}$$

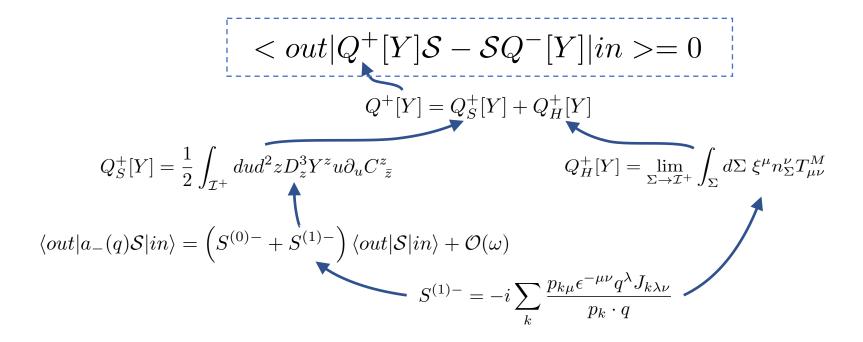
 $C_{zz}|_{\mathcal{J}^+_-} = C_{zz}|_{\mathcal{J}^-_+}, \quad m_B|_{\mathcal{J}^+_-} = m_B|_{\mathcal{J}^-_+},$

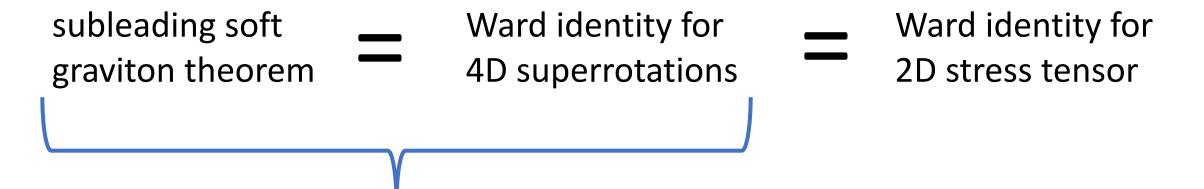
radiative
$$h_{\mu\nu}(x) = \sum_{\alpha=\pm} \int \frac{d^3q}{(2\pi)^3} \frac{1}{2\omega_q} \left[\epsilon_{\mu\nu}^{\alpha*}(\vec{q}) a_\alpha(\vec{q}) e^{iq\cdot x} + \epsilon_{\mu\nu}^\alpha(\vec{q}) a_\alpha(\vec{q})^\dagger e^{-iq\cdot x} \right]$$

$$C_{\bar{z}\bar{z}} = 2 \lim_{r \to \infty} \frac{1}{r} \partial_{\bar{z}} x^\mu \partial_{\bar{z}} x^\nu \sum_{\alpha=\pm} \int \frac{d^3q}{(2\pi)^3} \frac{1}{2\omega_q} \left[\epsilon_{\mu\nu}^{\alpha*}(\vec{q}) a_\alpha(\vec{q}) e^{-i\omega_q u - i\omega_q r(1 - \cos\theta)} + h.c. \right]$$
saddle
$$C_{\bar{z}\bar{z}} = -\frac{i}{4\pi^2} \hat{\epsilon}_{\bar{z}\bar{z}}^+ \int_0^\infty d\omega_q [a_-(\omega_q \hat{x}) e^{-i\omega_q u} - a_+(\omega_q \hat{x})^\dagger e^{i\omega_q u}]$$

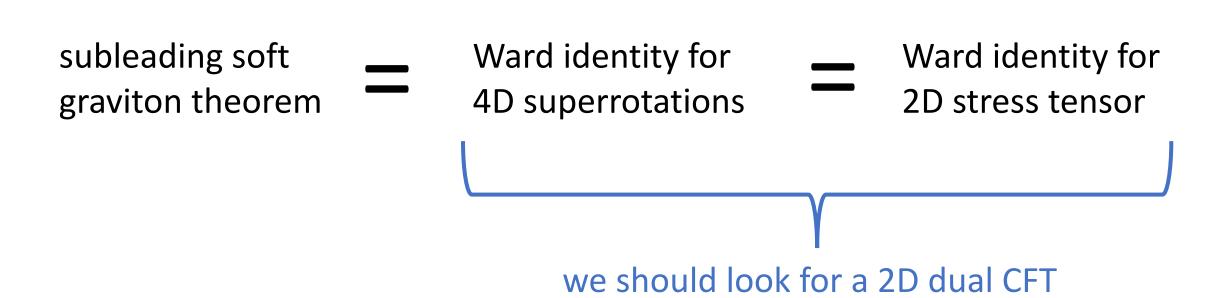
$$\langle out|a_{-}(q)\mathcal{S}|in\rangle = \left(S^{(0)-} + S^{(1)-}\right)\langle out|\mathcal{S}|in\rangle + \mathcal{O}(\omega)$$
$$S^{(0)-} = \sum_{k} \frac{(p_k \cdot \epsilon^{-})^2}{p_k \cdot q} \qquad S^{(1)-} = -i\sum_{k} \frac{p_{k\mu}\epsilon^{-\mu\nu}q^{\lambda}J_{k\lambda\nu}}{p_k \cdot q}$$

Soft Thm = Ward Id

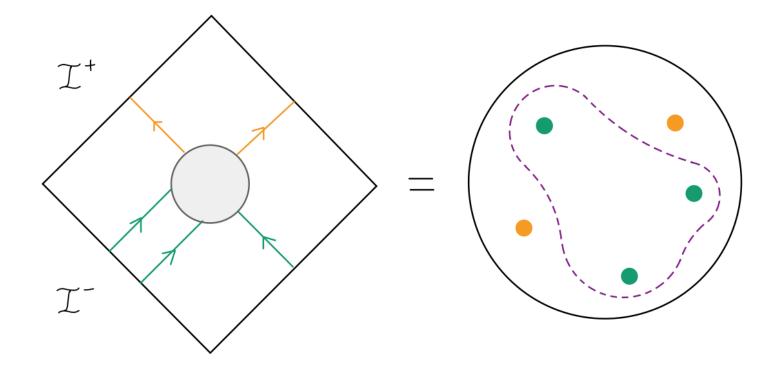




the asymptotic symmetry is physical

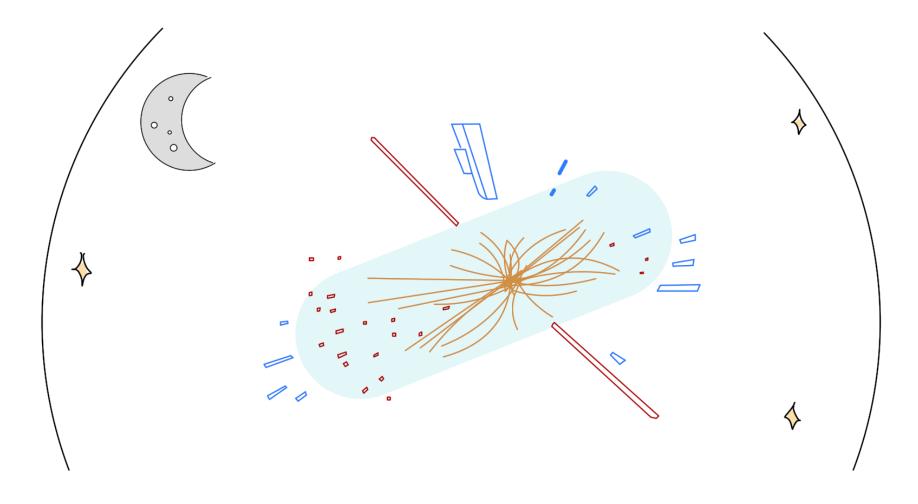


4D Soft Mode = 2D Current



The Celestial Conjecture:

scattering in asymptotically flat spacetimes is dual to a CFT living on the celestial sphere



4D Amplitude = 2D Correlator

4D Lorentz invariance = 2D global conformal symmetry

$$\langle \mathfrak{G}_{\Delta_1}^{\pm}(z_1, \bar{z}_1) ... \mathfrak{G}_{\Delta_n}^{\pm}(z_n, \bar{z}_n) \rangle = \prod_{i=1}^n \int_0^\infty d\omega_i \omega_i^{\Delta_i - 1} \langle out | \mathscr{S} | in \rangle$$

If we go to a boost basis, amplitudes transform as CFT correlators under the Lorentz group.



$$\underset{\text{scalar}}{\text{massive}} \quad \widetilde{\mathcal{A}}(\Delta_i, \vec{w_i}) \equiv \prod_{k=1}^n \int_{H_{d+1}} [d\hat{p}_k] \, G_{\Delta_k}(\hat{p}_k; \vec{w}_k) \ \mathcal{A}(\pm m_i \hat{p}_i^{\mu})$$

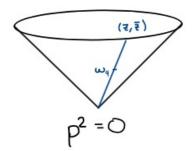
$$\boldsymbol{m} = \boldsymbol{0} \qquad \widetilde{\mathcal{A}}(\Delta_i, \vec{w}_i) \equiv \prod_{k=1}^n \int_0^\infty d\omega_k \omega_k^{\Delta - 1} \,\mathcal{A}(\pm \omega_k q_k^{\mu})$$

Lorentz covariance guaranteed by this choice of wavepackets, with u-direction captured by a continuous spectrum

 $\Delta = 1 + i\lambda$

Meanwhile translations shift the weight

$$p^{\mu} = q^{\mu} e^{\partial_{\Delta}} \Leftrightarrow \Delta \mapsto \Delta + 1$$

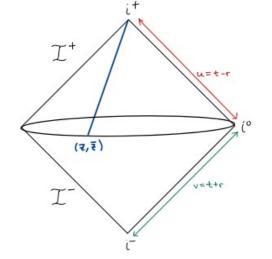


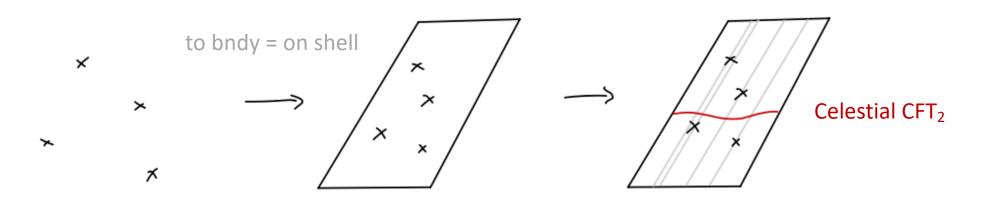
4D Amplitude = 4D Correlator

LSZ \Leftrightarrow Extrapolate Dict.

$$\langle out|S|in\rangle_{boost} = \prod_{i} \lim_{r \to \infty} \int_{-\infty}^{\infty} \mathrm{d}\nu_{i} \,\nu_{i}^{-\Delta_{i}} \,\langle r\Phi(\nu_{1}, r, z_{1}, \bar{z}_{1})...r\Phi(\nu_{n}, r, z_{n}, \bar{z}_{n})\rangle$$

$$\nu = \{u, v\}$$





perturbative bulk

Carrollian CFT₃

Operator Spectrum

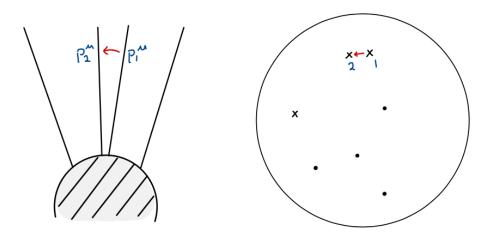
Fock Space \leftrightarrow 4D Hilbert Space \leftrightarrow 2D States \leftrightarrow 2D Operators

$$\mathfrak{G}_{\Delta}(z,\bar{z}) \equiv \int_{-\infty}^{\infty} \mathrm{d}u \, u^{-\Delta} \lim_{r \to \infty} \left[r^{\delta} \Phi(u,r,z,\bar{z}) \right]$$

$$: \mathcal{O}^{(\rho)}\mathcal{O}:_{\Delta}(z,\bar{z}) \; \equiv \; \int_{0}^{\infty} d\omega \, \omega^{\Delta-\rho-1} \, \int_{0}^{\omega} \, d\omega_1 \, \omega_1^{\rho-1} \, a^{\dagger}(\omega_1,z,\bar{z}) a^{\dagger}(\omega-\omega_1,z,\bar{z})$$

Collinear Limit = OPE

$$\begin{split} \mathfrak{G}_{\Delta_{1},+2}(z_{1},\bar{z}_{1})\mathfrak{G}_{\Delta_{2},+2}(z_{2},\bar{z}_{2}) &\sim -\frac{\kappa}{2}\frac{\bar{z}_{12}}{z_{12}}B(\Delta_{1}-1,\Delta_{2}-1)\mathfrak{G}_{\Delta_{1}+\Delta_{2},+2}(z_{2},\bar{z}_{2})+\dots, \\ \mathfrak{G}_{\Delta_{1},+2}(z_{1},\bar{z}_{1})\mathfrak{G}_{\Delta_{2},-2}(z_{2},\bar{z}_{2}) &\sim -\frac{\kappa}{2}\frac{\bar{z}_{12}}{z_{12}}B(\Delta_{1}-1,\Delta_{2}+3)\mathfrak{G}_{\Delta_{1}+\Delta_{2},-2}(z_{2},\bar{z}_{2}) \\ &- \frac{\kappa}{2}\frac{z_{12}}{\bar{z}_{12}}B(\Delta_{1}+3,\Delta_{2}-1)\mathfrak{G}_{\Delta_{1}+\Delta_{2},+2}(z_{2},\bar{z}_{2})+\dots, \end{split}$$

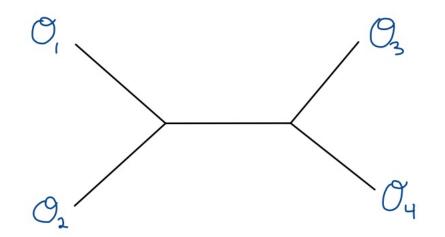


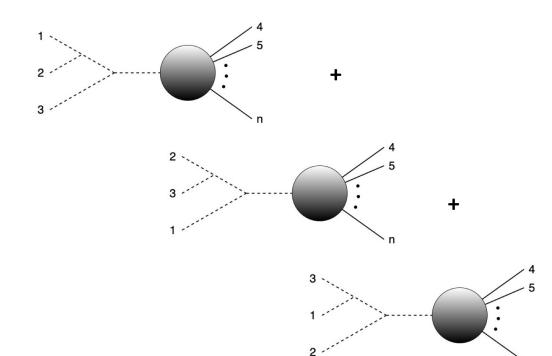
Celestial OPE

Celestial OPE









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Celestial OPE

Celestial OPE vs Feynman Diagrams

$$\mathcal{O}_{\Delta_{1}}(z_{1},\bar{z}_{1})\mathcal{O}_{\Delta_{2}}(z_{2},\bar{z}_{2})\mathcal{O}_{\Delta_{3}}(z_{3},\bar{z}_{3}) \sim \left(\frac{1}{z_{13}z_{23}\bar{z}_{13}\bar{z}_{23}}\mathcal{C}_{1} + \frac{(z_{23}\bar{z}_{23})^{\Delta_{1}-2}}{(z_{13}\bar{z}_{13})^{\Delta_{1}}}\mathcal{C}_{2}\right)\mathcal{O}_{\Delta_{1}+\Delta_{2}+\Delta_{3}-4}(z_{3},\bar{z}_{3})$$

$$\mathcal{O}_{\Delta_{2}}(z_{2},\bar{z}_{2})\mathcal{O}_{\Delta_{3}}(z_{3},\bar{z}_{3}) \sim \int d\Delta \frac{C_{\Delta_{2},\Delta_{3}}}{(z_{23}\bar{z}_{23})^{\frac{1}{2}(\Delta_{2}+\Delta_{3}-\Delta)}}\mathcal{O}_{\Delta}(z_{3},\bar{z}_{3})$$

$$+ \int d\Delta d\sigma \frac{C_{\Delta_{2},\Delta_{3}}}{(z_{23}\bar{z}_{23})^{\frac{1}{2}(\Delta_{2}+\Delta_{3}-\Delta)}}\mathcal{R}_{\Delta}^{\sigma}(z_{3},\bar{z}_{3})$$

2D Radial Quantization → More Symmetries

For special weights, the SL(2,C) multiplets have primary descendants.

$$H^{k}(z, \bar{z}) := \lim_{\epsilon \to 0} \epsilon \, \mathfrak{G}_{k+\epsilon,2}(z, \bar{z}), \quad \Delta = k = 2, 1, 0, -1, \dots$$

Assuming these multiplets shorten, we have

$$H^{k}(z,\bar{z}) = \sum_{m=\frac{k-2}{2}}^{\frac{2-k}{2}} \bar{z}^{-\frac{k-2}{2}-m} H^{k}_{m}(z) , \qquad \qquad w^{p}_{n} = \frac{1}{\kappa} (p-n-1)! (p+n-1)! H^{-2p+4}_{n}(z) + \frac{1}{\kappa} (p-n-1)! (p+n-1)! (p+n-1)! H^{-2p+4}_{n}(z) + \frac{1}{\kappa} (p-n-1)! (p+n-1)! (p+n-1$$

2D Radial Quantization → More Symmetries

Complexifying the celestial sphere variables and defining a holomorphic commutator

$$[A,B](z) = \frac{1}{2\pi i} \oint_z dw A(w)B(z)$$

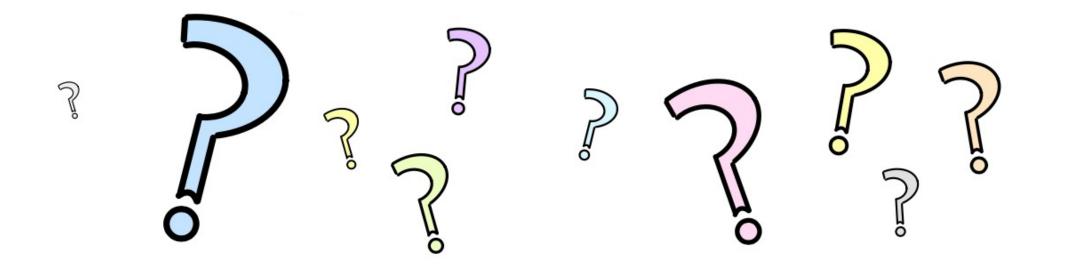
gives a $Lw_{1+\infty}$ symmetry algebra for appropriately rescaled modes

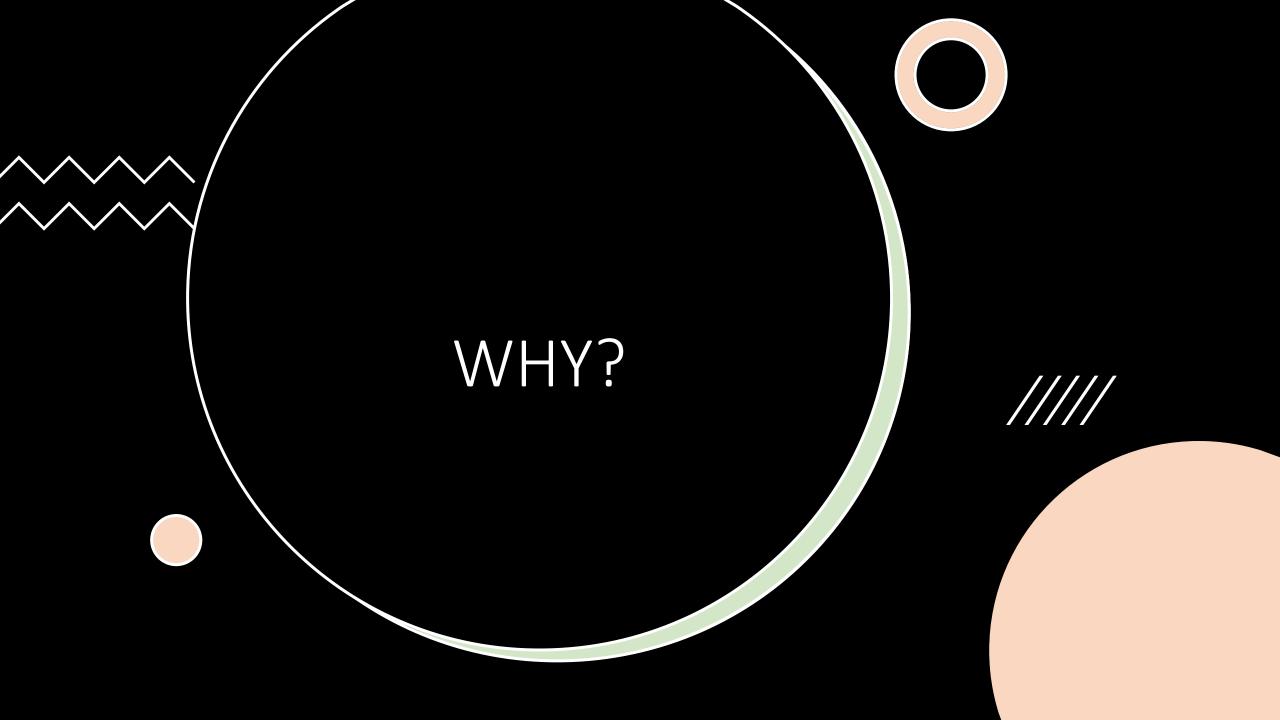
$$\left[w_{n}^{p}, w_{m}^{q}\right](z) = \left[n(q-1) - m(p-1)\right] w_{m+n}^{p+q-2}(z)$$

Do these symmetries beyond tree level, or the self-dual sector?

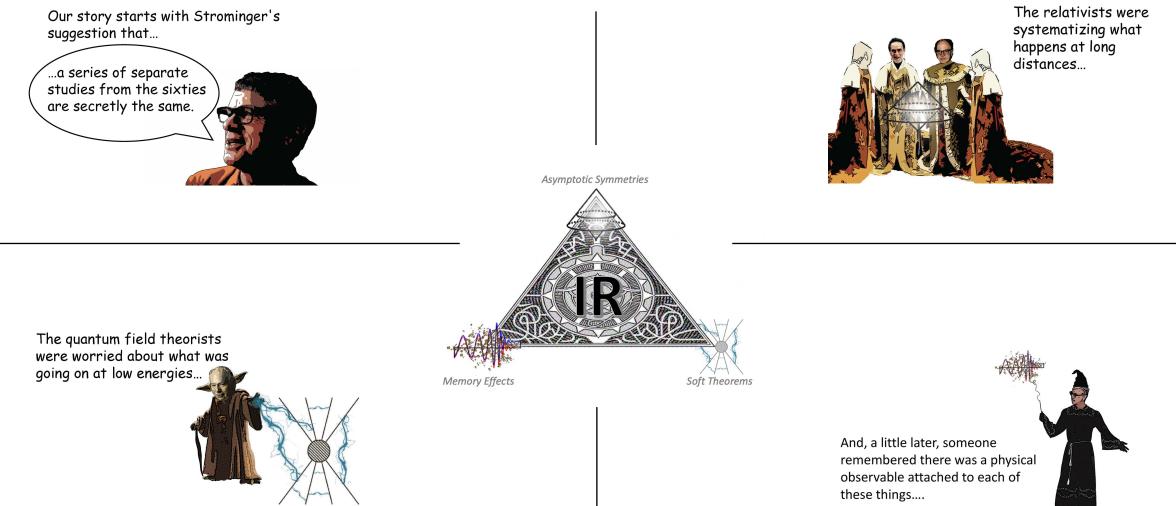
Can we realize them in the matter sector?

Can we really complexify the celestial sphere to define these currents?

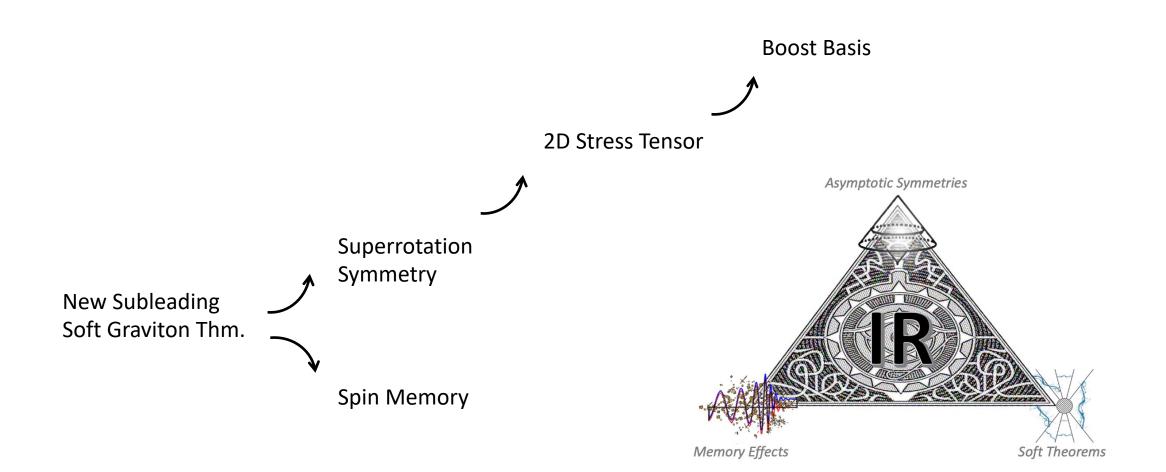




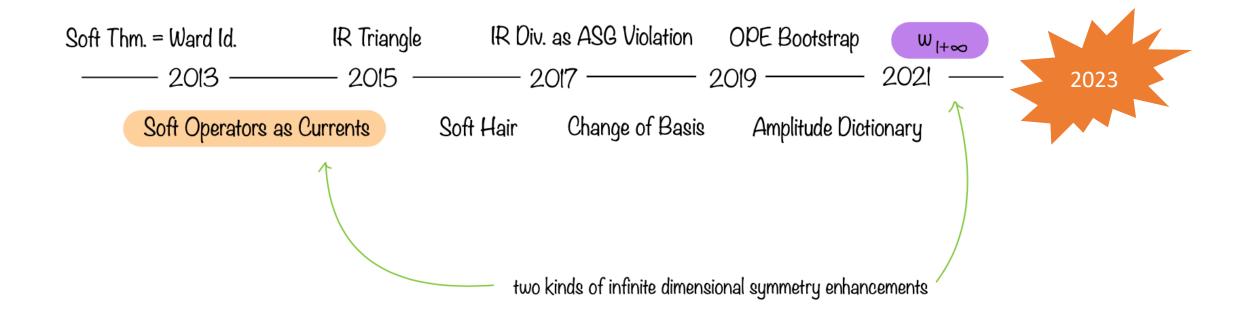
A Collision of Fields: Then

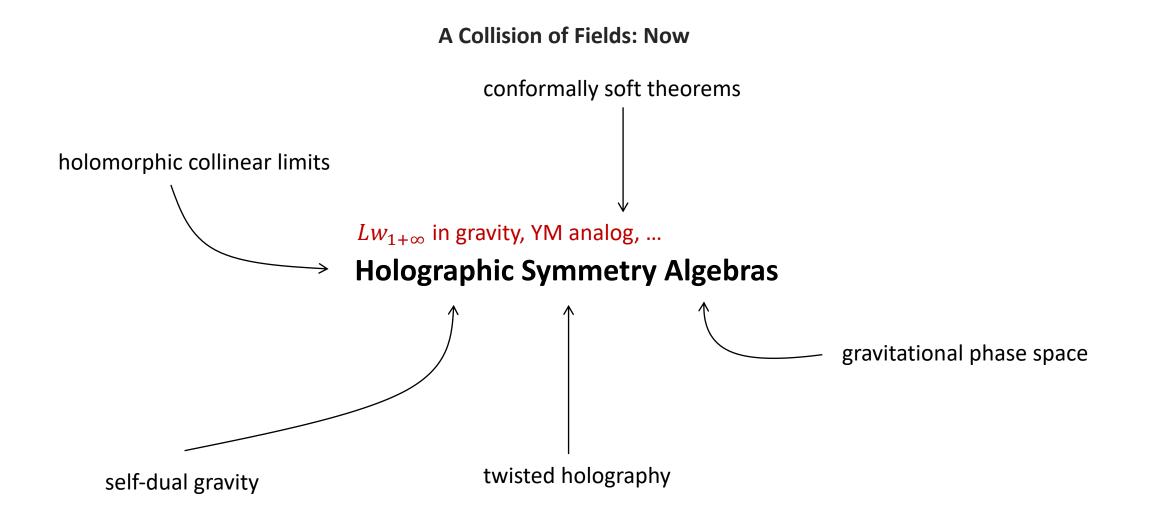


A Collision of Fields: Then

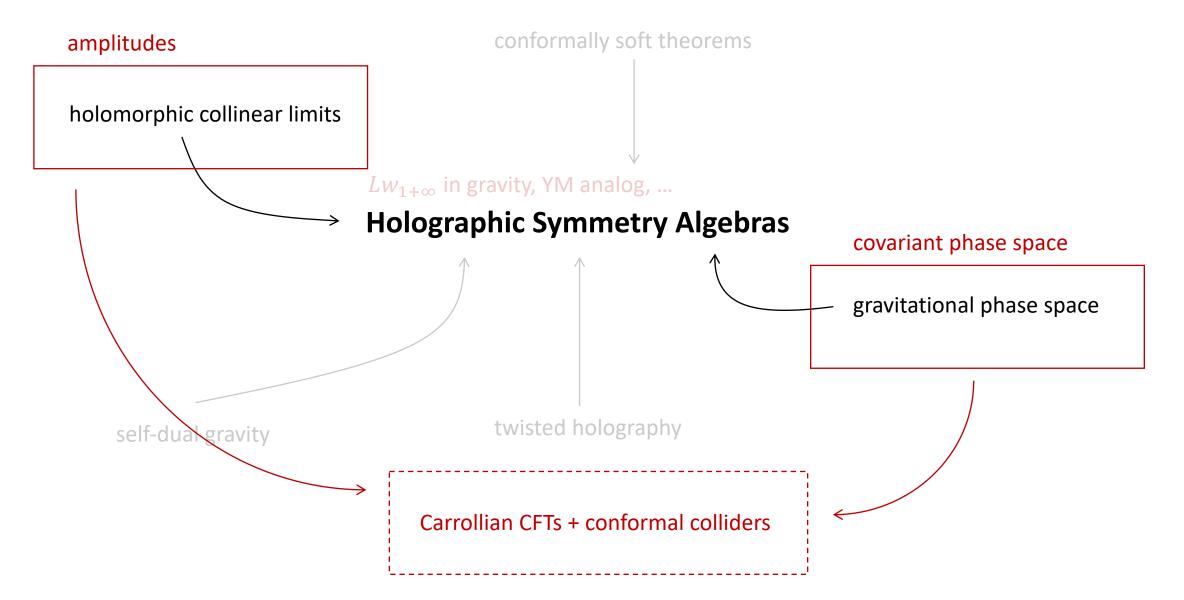


A Collision of Fields: Now

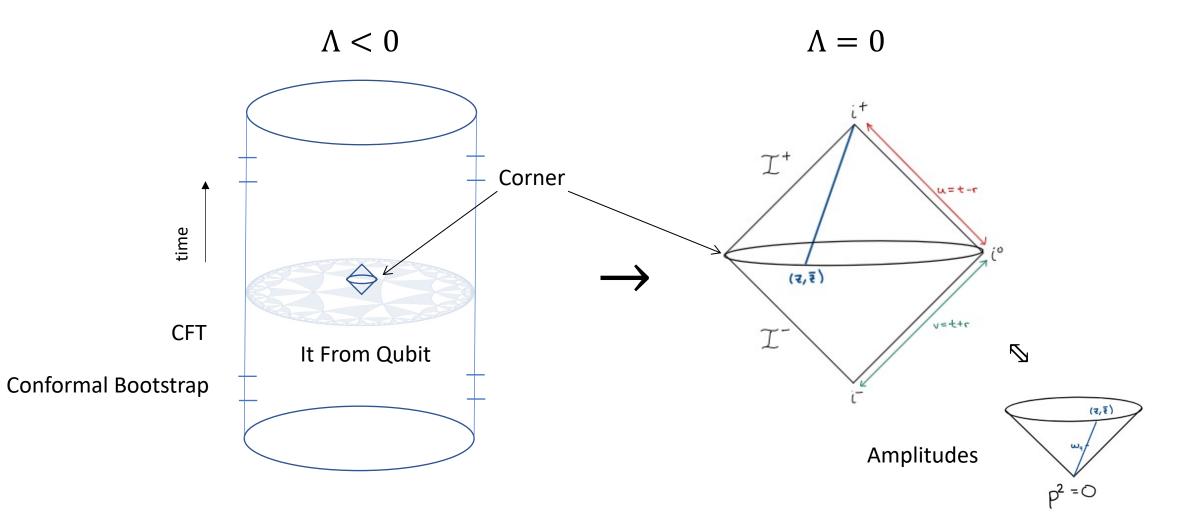




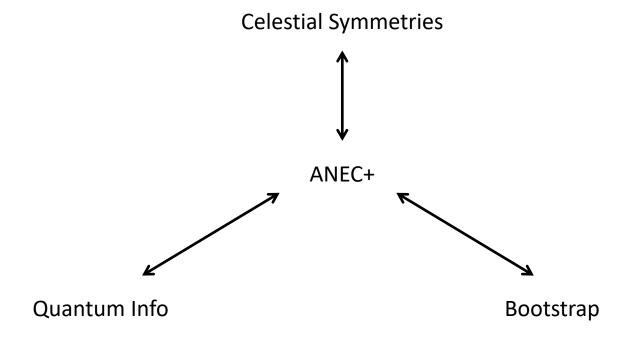
A Collision of Fields: Now

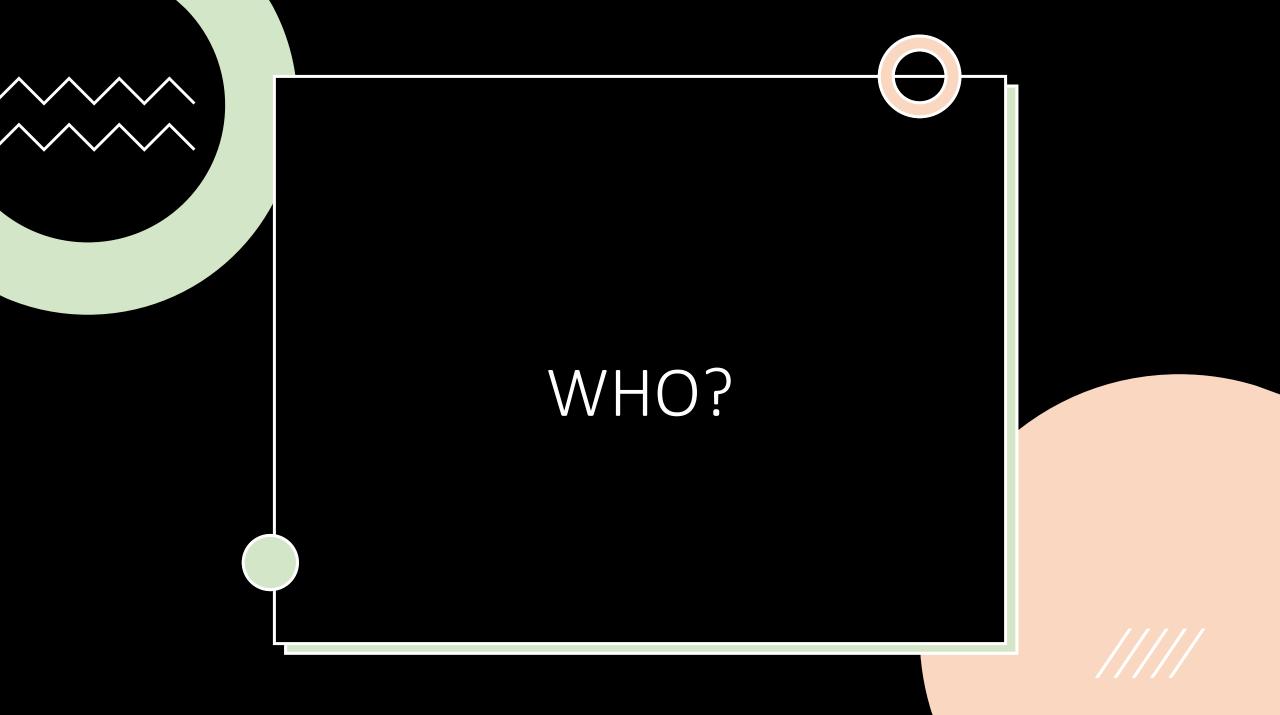


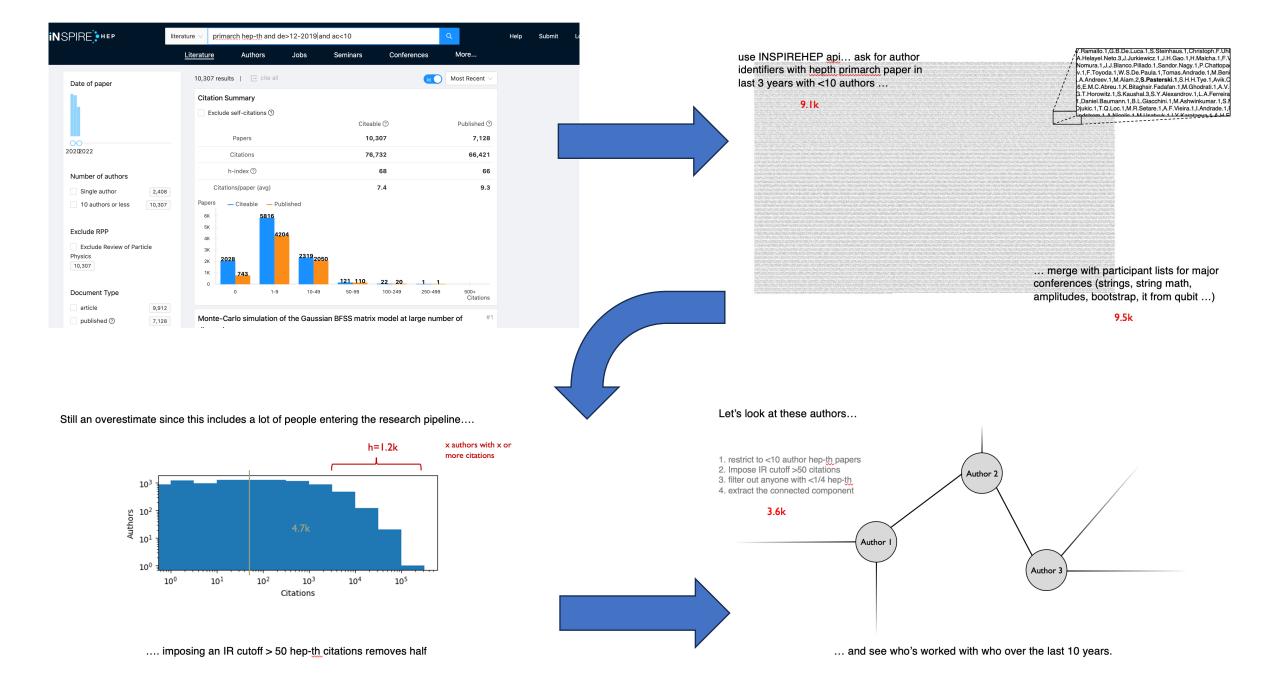
A Collision of Fields: Soon

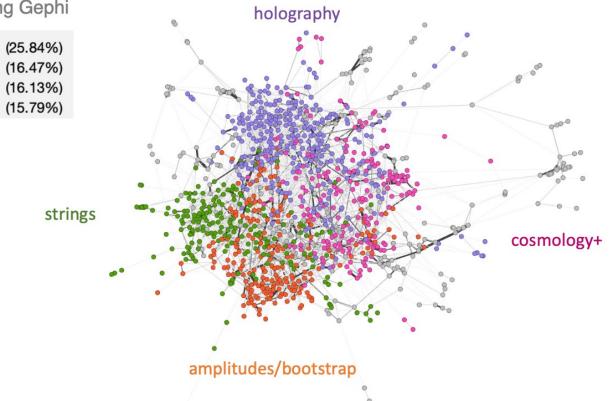


A Collision of Fields: Soon









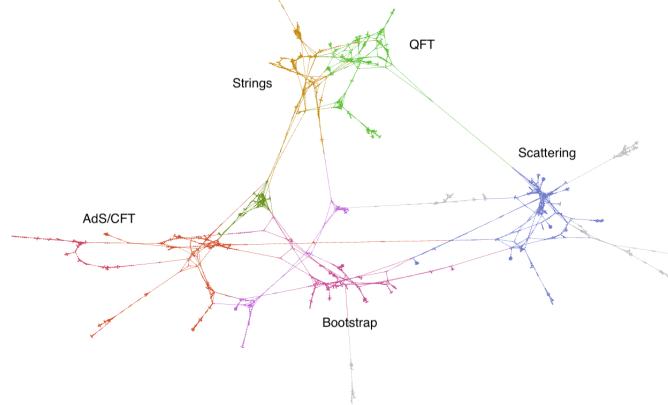
can identify communities using Gephi

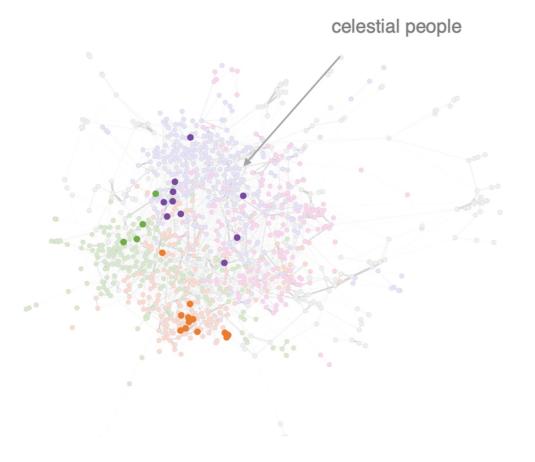
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| | 1 0 1 0 1 0 1 0 0 0 | 1 [[2617022], [2170581, 2122418, 2103461, 1866164, 1860825, 1 0 [[2058015], [1989007]] 1 [[2168963, 2063426, 2020643, 1856580, 1841069], [1869123]] 0 1 1 [[2166821, 2036329, 1960285], [2024912]] 0 1 1 [[2166821, 2036329, 1960285], [2024912]] 0 1 1 [[2166823, 2054698, 2016467, 1904870, 1870526], [1955845, 11 0 [[2178053, 2513318, 2162003, 2108965, 2028982, 2619683, 193 1 [[2610588, 2514138, 2120761, 2072456, 2072539, 2058002, 193 | 1 [[2617022], [2170581, 2122418, 2103461, 1866164, 1860825, 11 ['derivative scalar theory', 'bp 0 [[2058015], [1989007]] ['massless schwinger model', 1 [[2168963, 2063426, 2020643, 1856580, 1841069], [1869123]] ['scattering radiation reaction' 0 1 [[2168821, 2036329, 1960285], [2024912]] ['complexity equals', 'extrema 0 1 [[2166821, 2036329, 1960285], [2024912]] ['complexity equals', 'extrema 0 [[2616328, 2054698, 2016467, 1904870, 1870526], [1955845, 11 ['einstein cartan gravity', 'finito' 0 [[2178053, 2513318, 2162003, 2108965, 2028982, 2619683, 19; ['gravitational chiral anomaly', 1 [[2610588, 2514138, 2120761, 2072456, 2072539, 2058002, 19i ['super yang mills', 'point intego | |

themes. feel free to use your knowledge of the keywords

Certainly. Here are the ten main themes distilled from the provided keywords:

- 1. Holography and AdS/CFT
- 2. Black Hole Physics
- 3. Conformal and Quantum Field Theories
- 4. String Theory and Dualities
- 5. Supersymmetry and Supergravity
- 6. Topological and Geometric Methods
- 7. Entanglement and Quantum Mechanics
- 8. Gauge Theories
- 9. Cosmology and the Universe
- 10. Quantum Gravity and General Relativity





can identify communities using Gephi

| 6 | (25.84%) |
|---|----------|
| 1 | (16.47%) |
| 0 | (16.13%) |
| 2 | (15.79%) |



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